

# Single-artificial-atom lasing and its suppression due to strong pumping

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in collaboration with

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A.M. Zagoskin  
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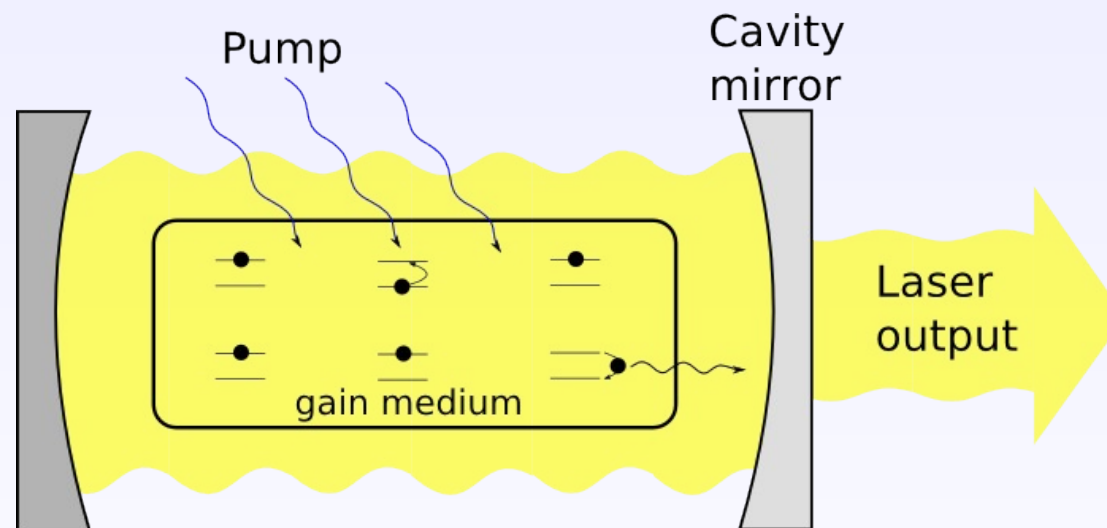
(RIKEN, Japan)



arXiv:0803.1209

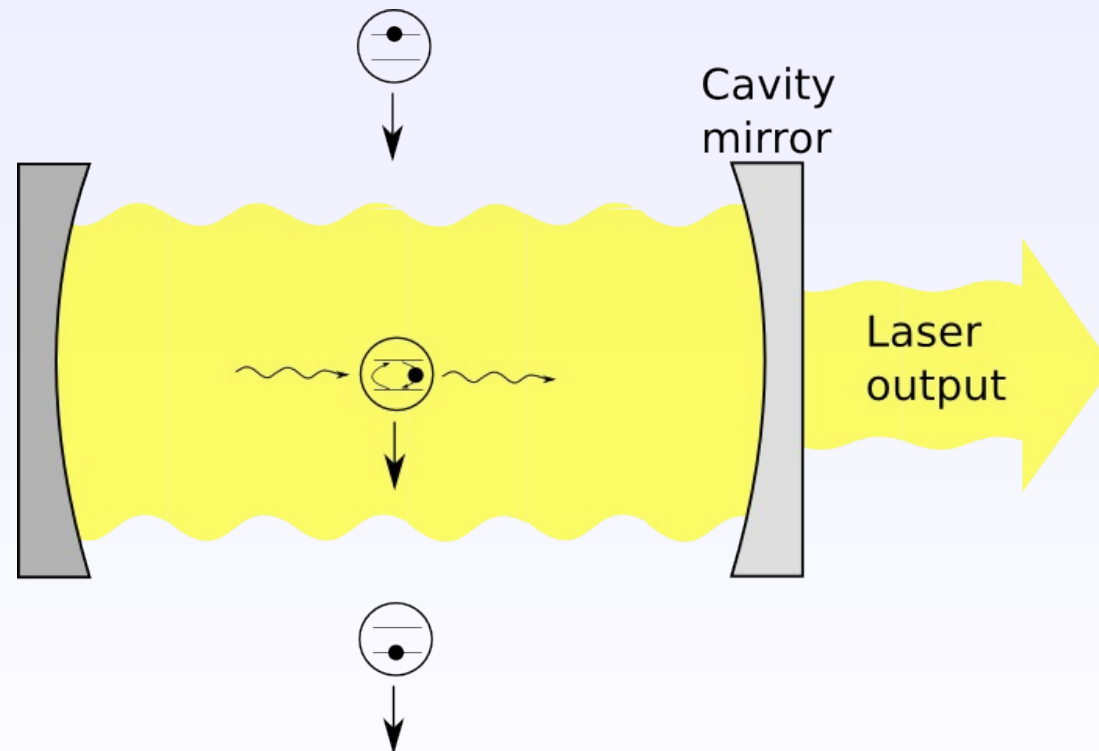
# Laser overview

- *A typical laser consist of*
  - *A cavity*
  - *A gain medium*
  - *A pump source*



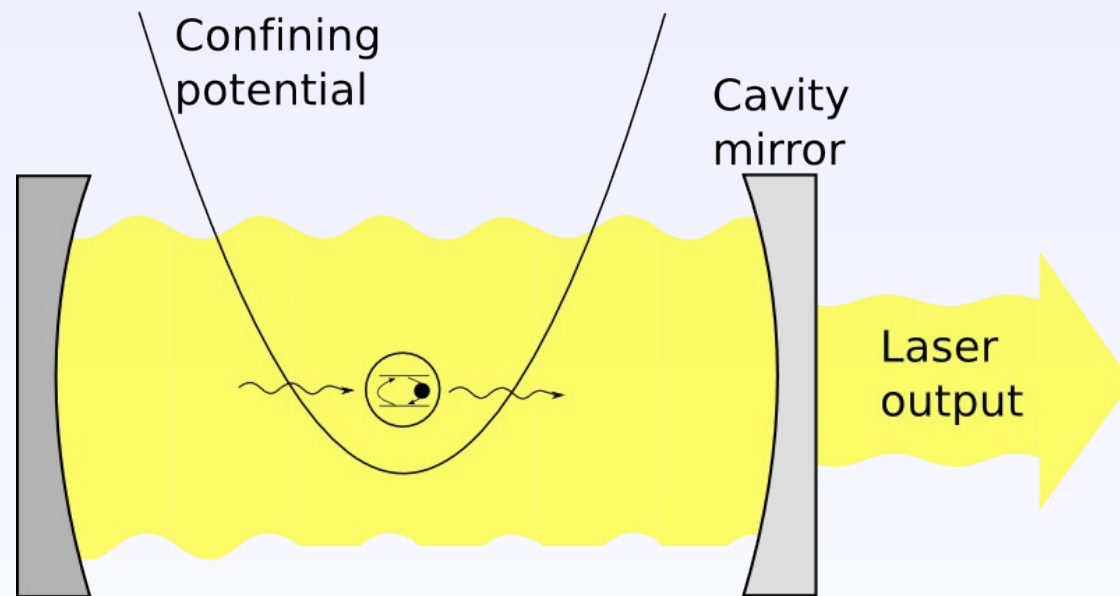
# Single-atom laser

- *The pump medium consist of a single atom (at the time)*
  - *First experiment used excited Rydberg atoms traveling through a cavity. Inside the cavity they interact with and transfer energy to the cavity field. [K. An et al. PRL **73**, 3375 (1994)]*



# Single-atom laser

- *The pump medium consist of a single atom*
  - *Another experiment used an atom (Cesium) in a trap inside the cavity [J. McKeever et al., Nature 245, 268 (2003)]*



# Requirements for single-atom lasing

- *Requirements to achieve single-atom lasing:*
  - *Strong atom-cavity coupling*

$$g \gg \kappa$$

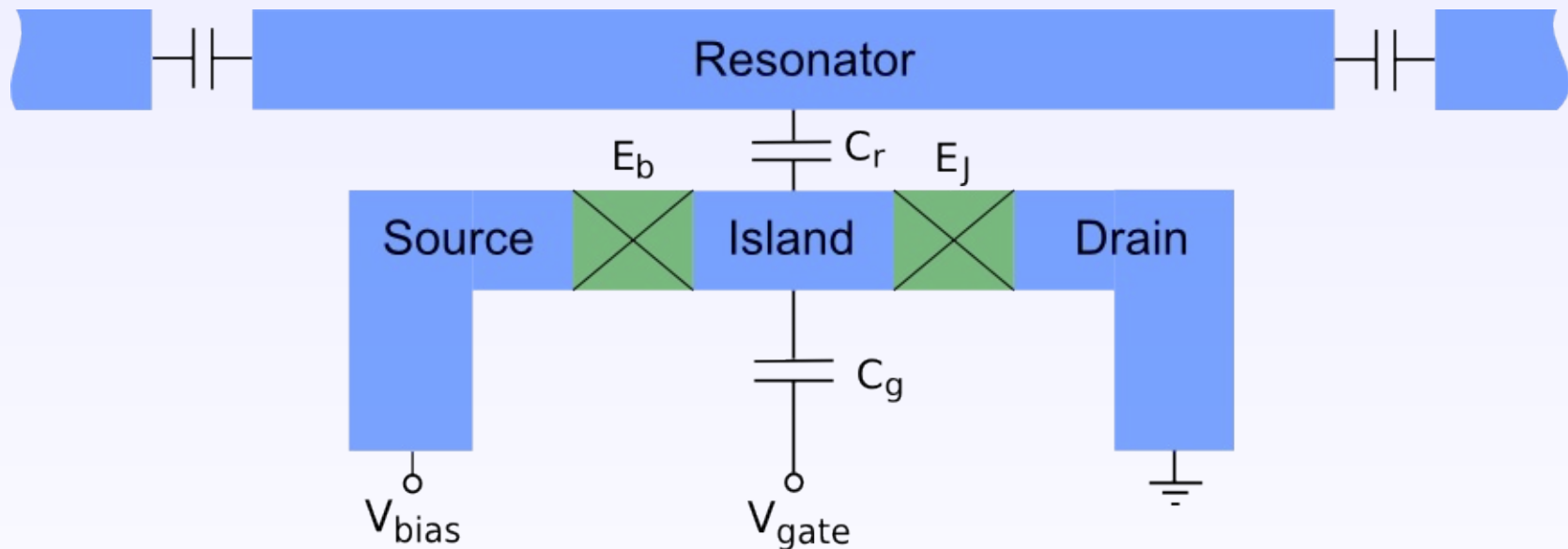
- *The life-time of the atom larger than interaction time*

$$g \gg \gamma$$

- *Both of these requirements are challenging in quantum optics*
  - *relatively small atom dipole moment*
  - *difficult to trap atoms in the cavity*

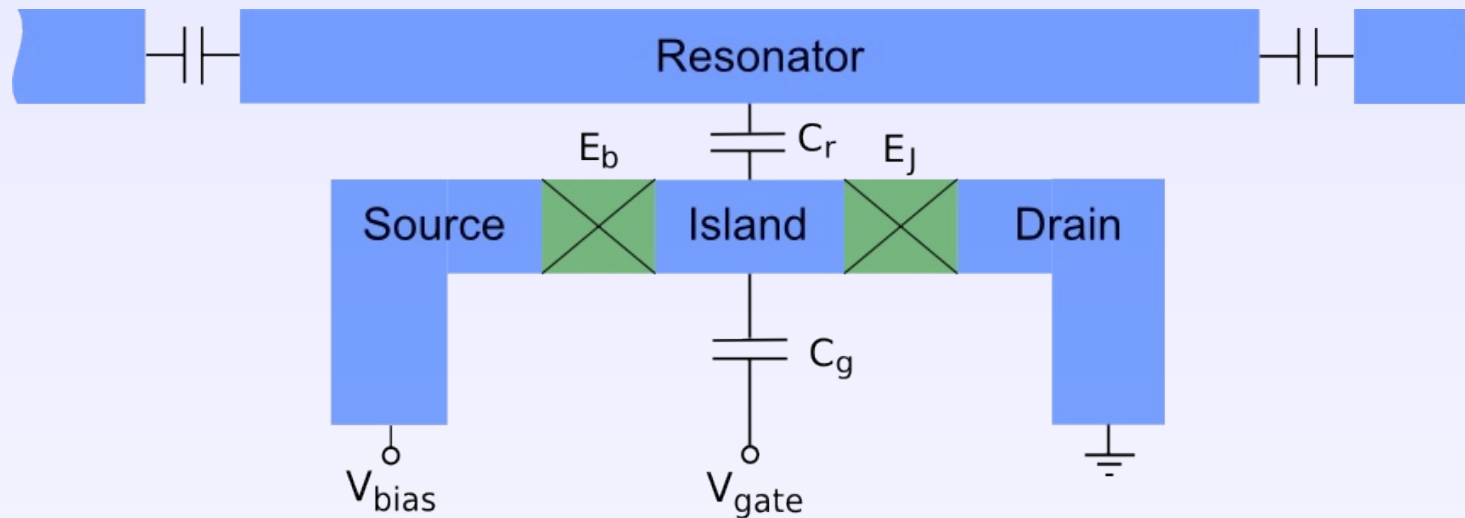
# Single-artificial-atom laser

- Recently a “single-atom laser” has been realized with a superconducting charge qubit (“artificial atom”) coupled to a transmission line resonator (“cavity”)
  - Astafiev et al., *Nature* **449**, 588 (2007)



# Single-artificial-atom laser: advantages

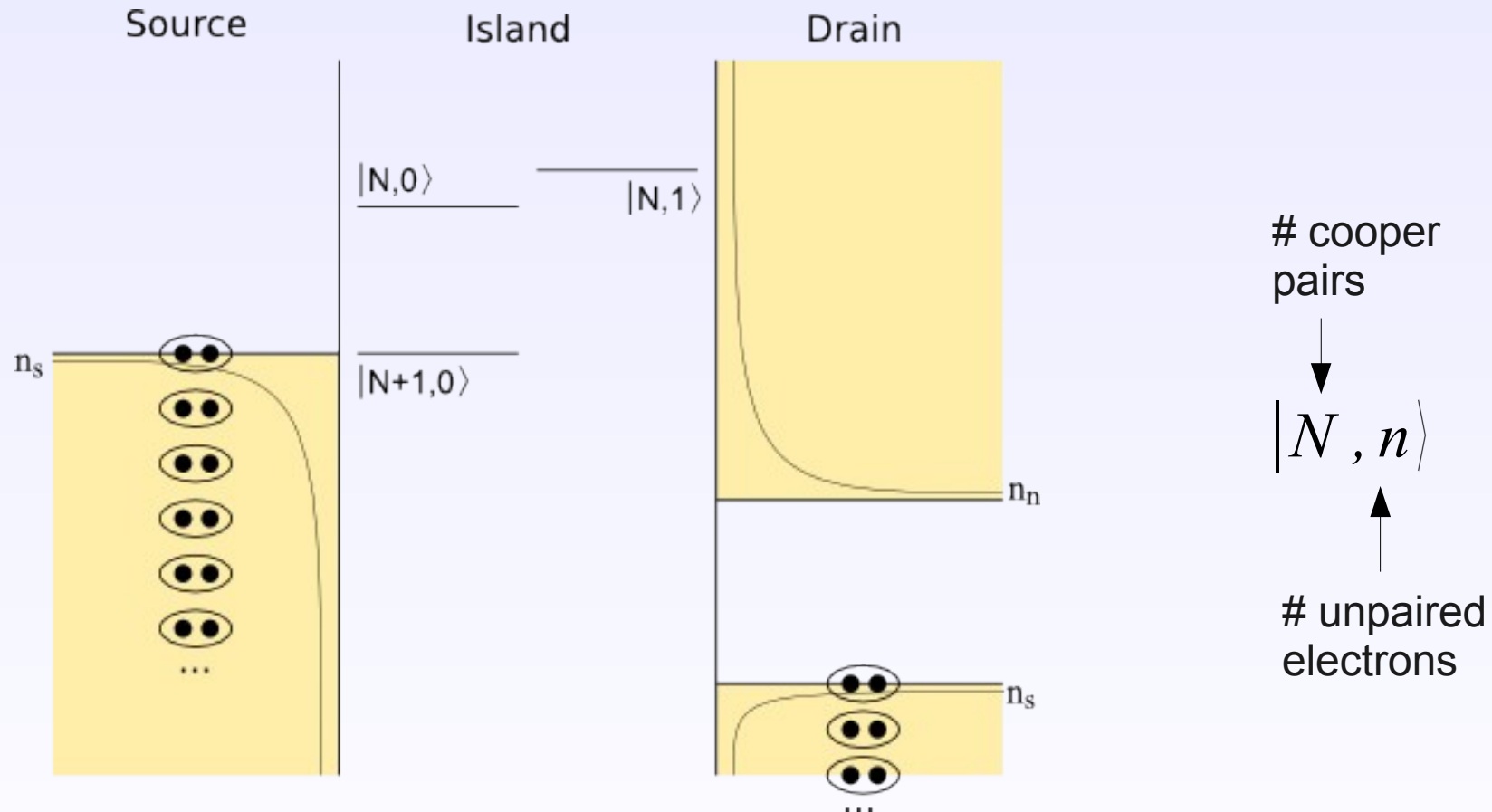
- *Astafiev et al., Nature* **449**, 588 (2007)



- *Advantages:*
  - *The artificial atom is “fixed inside the cavity”*
  - *Due to large dipole moment, strong atom-cavity coupling is possible*

# Single-artificial-atom laser: pumping cycle

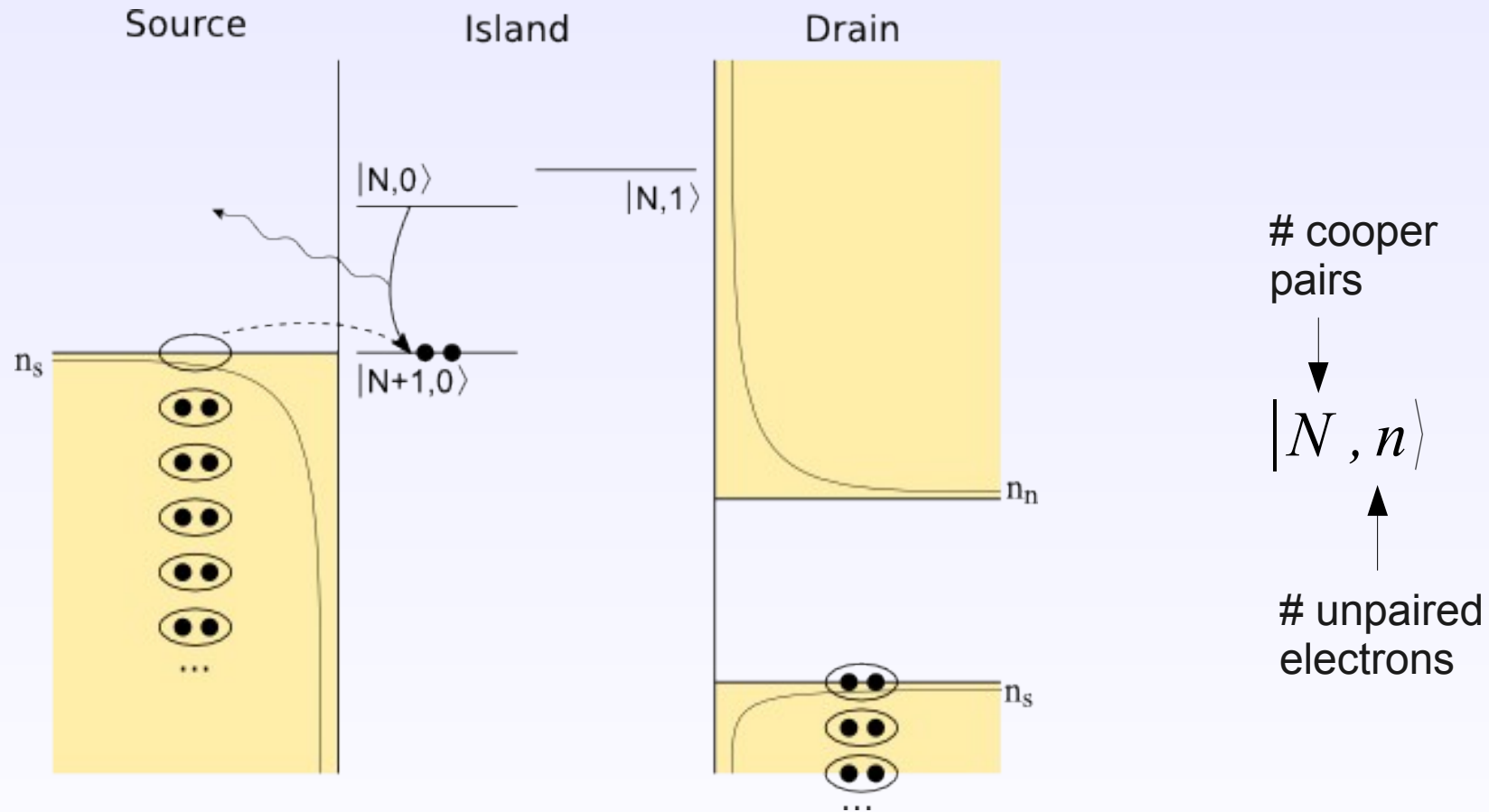
- *Population inversion of the artificial atom by a Josephson quasi-particle cycle:*





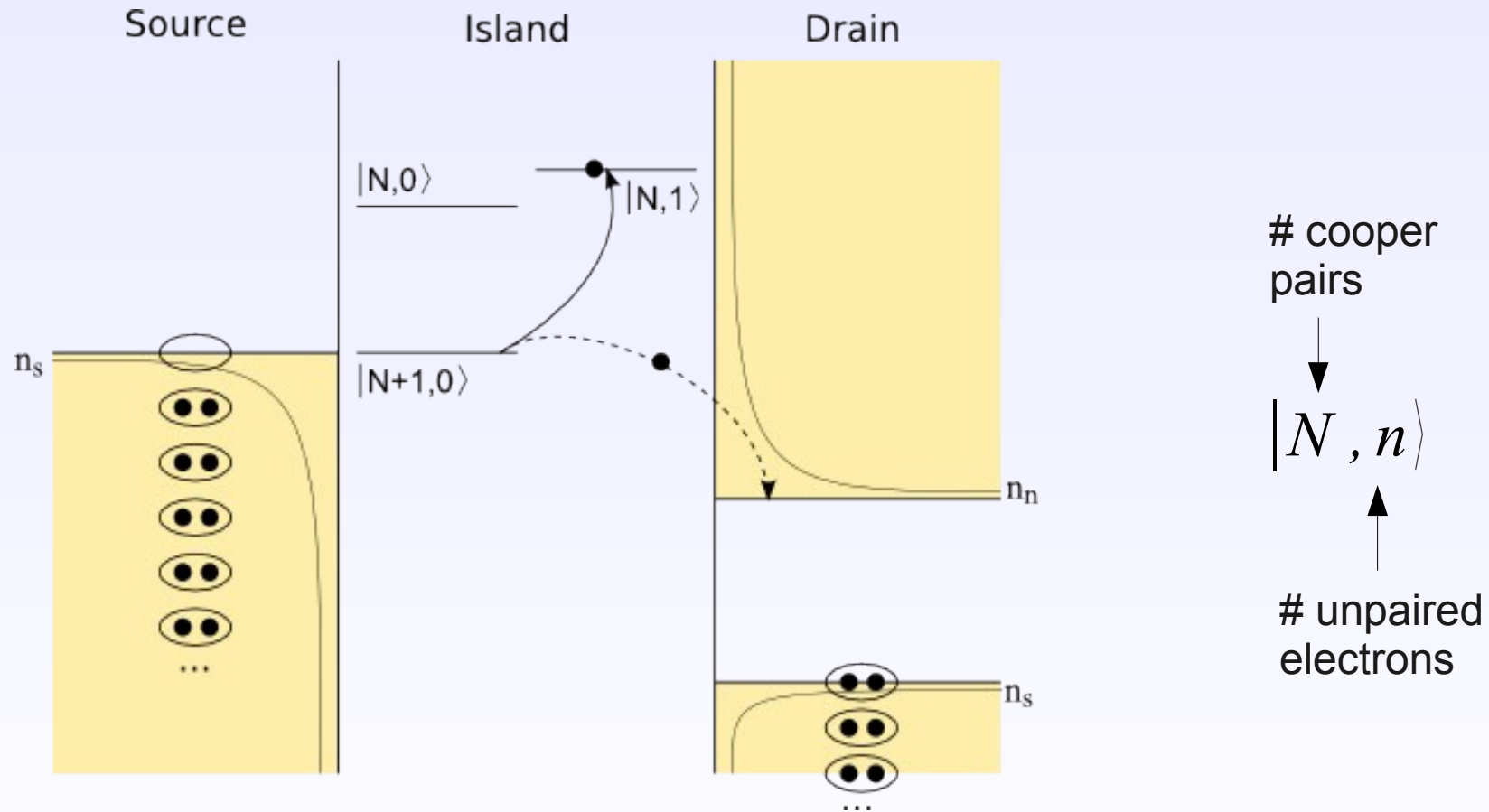
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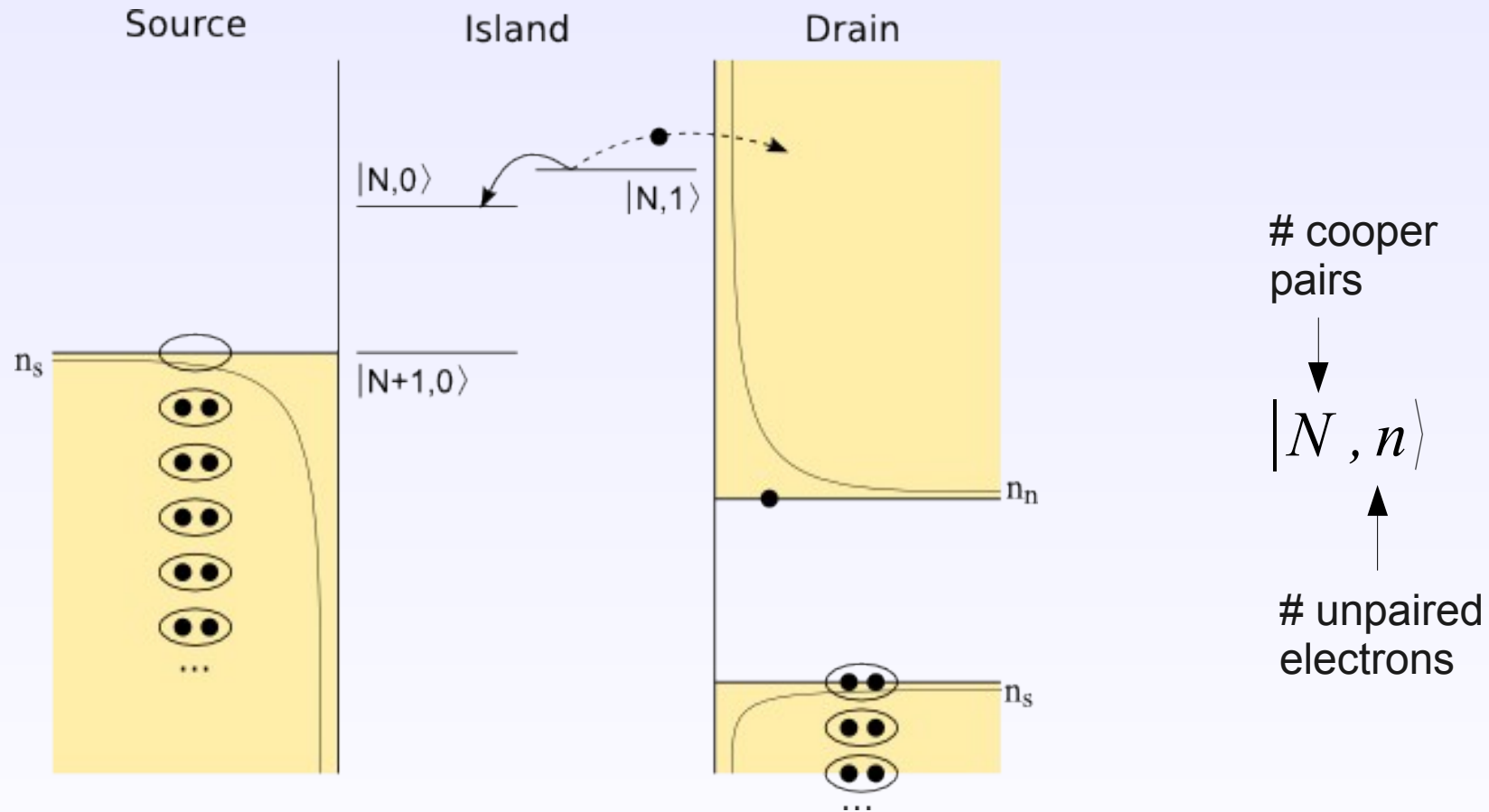
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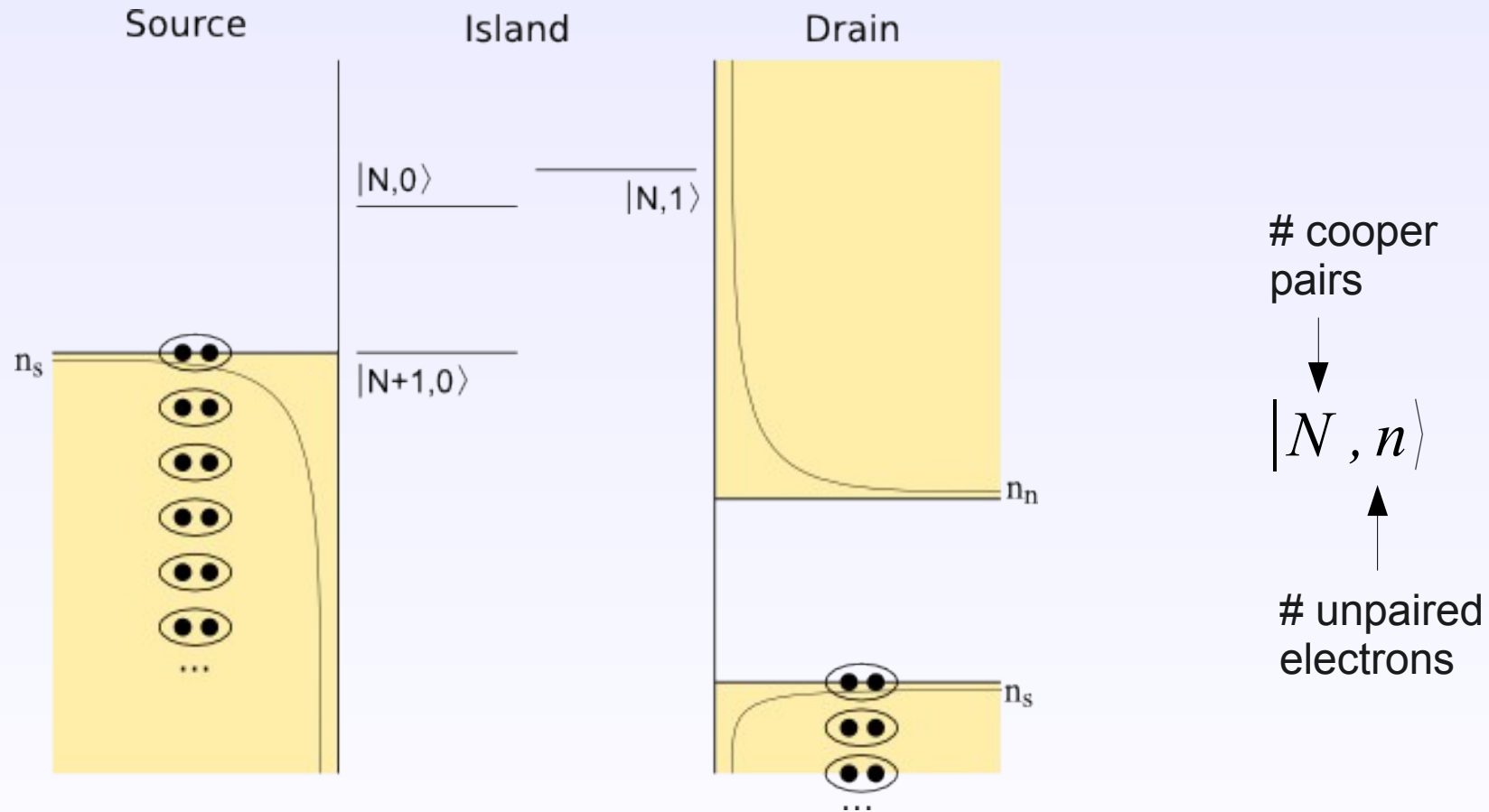
# Single-artificial-atom laser: pumping cycle

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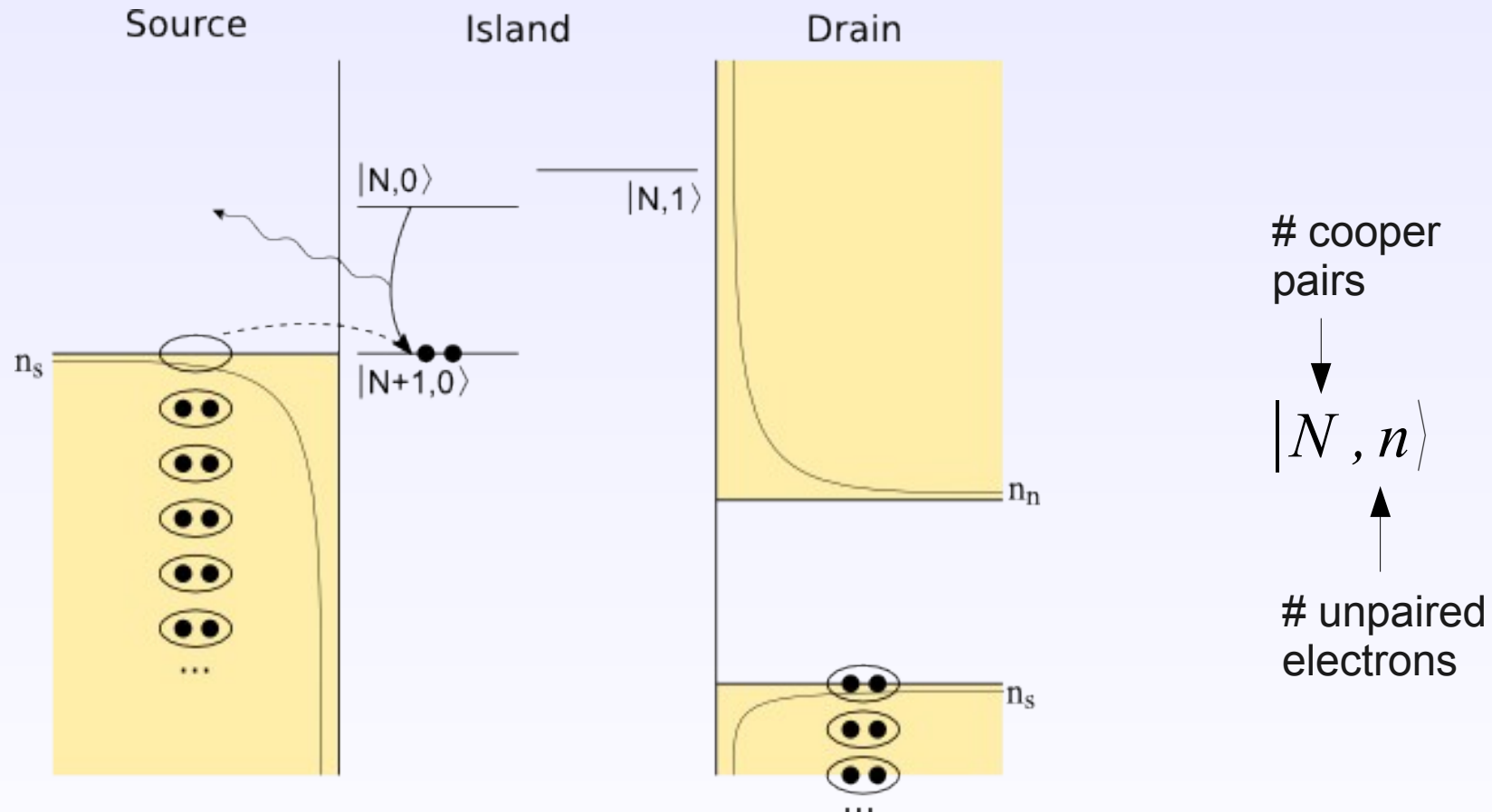
# Single-artificial-atom laser: pumping cycle

- *Population inversion of the artificial atom by a Josephson quasi-particle cycle:*



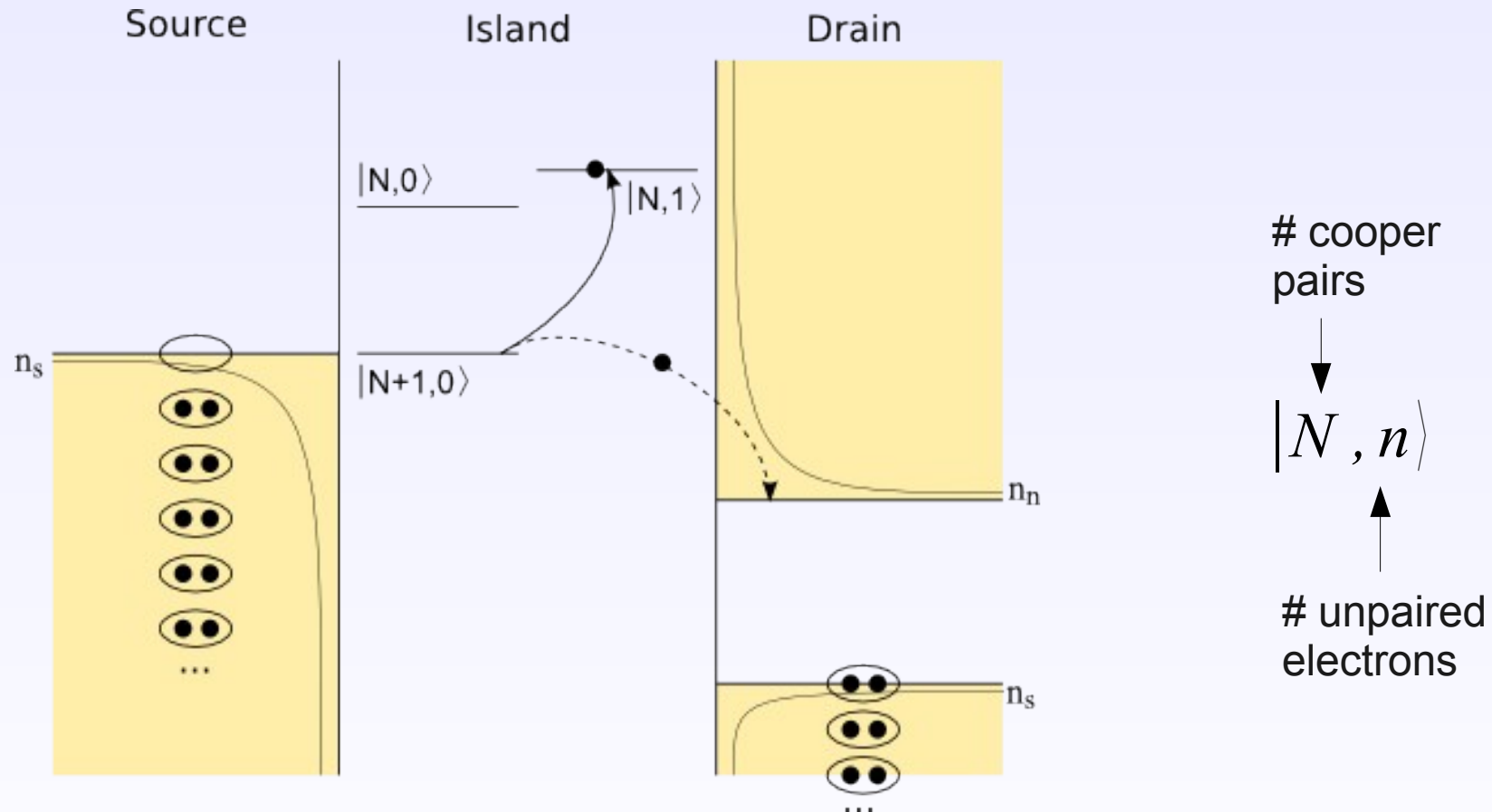
# Single-artificial-atom laser: pumping cycle

- *Population inversion of the artificial atom by a Josephson quasi-particle cycle:*



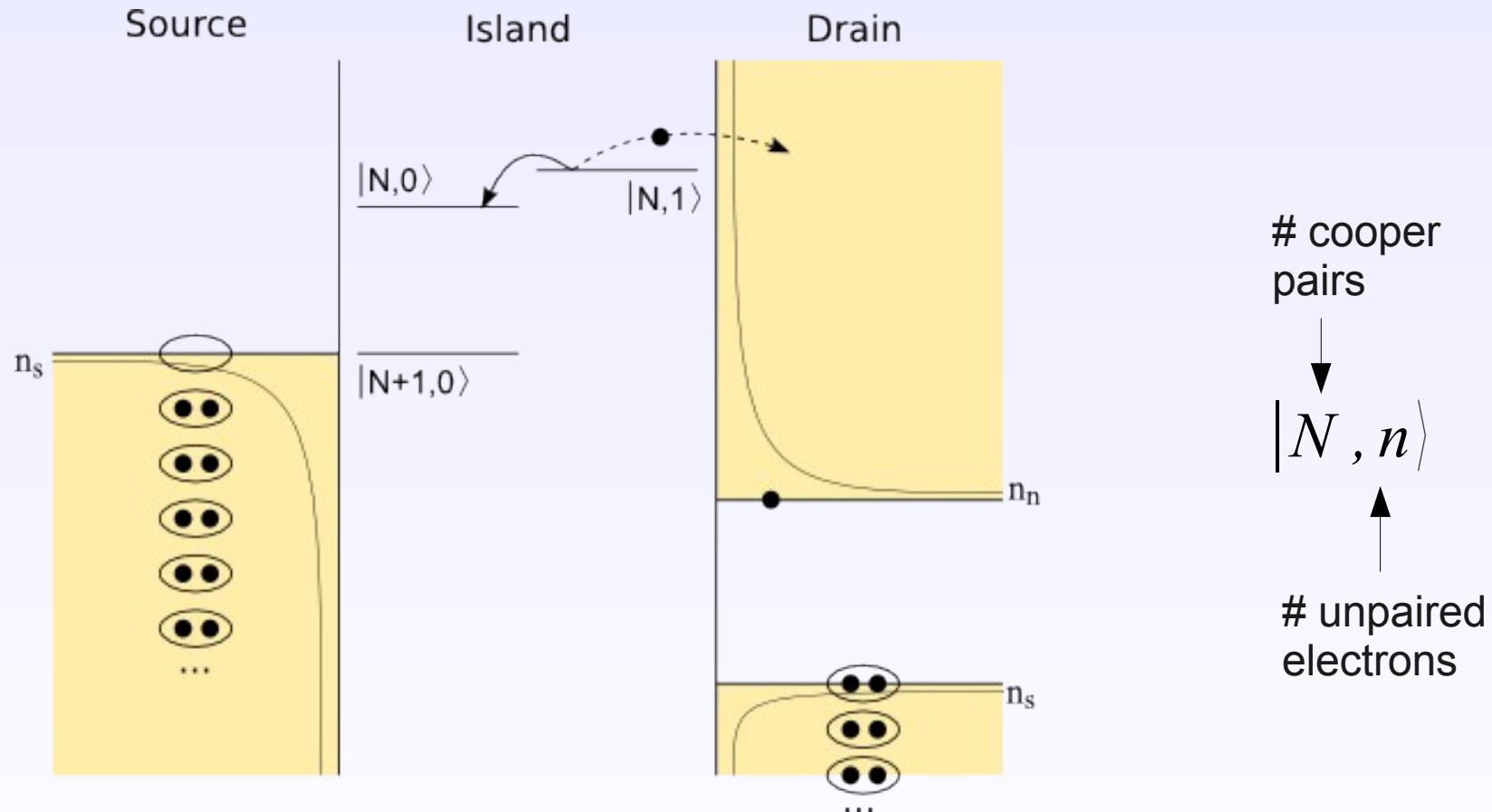
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- *Population inversion of the artificial atom by a Josephson quasi-particle cycle:*

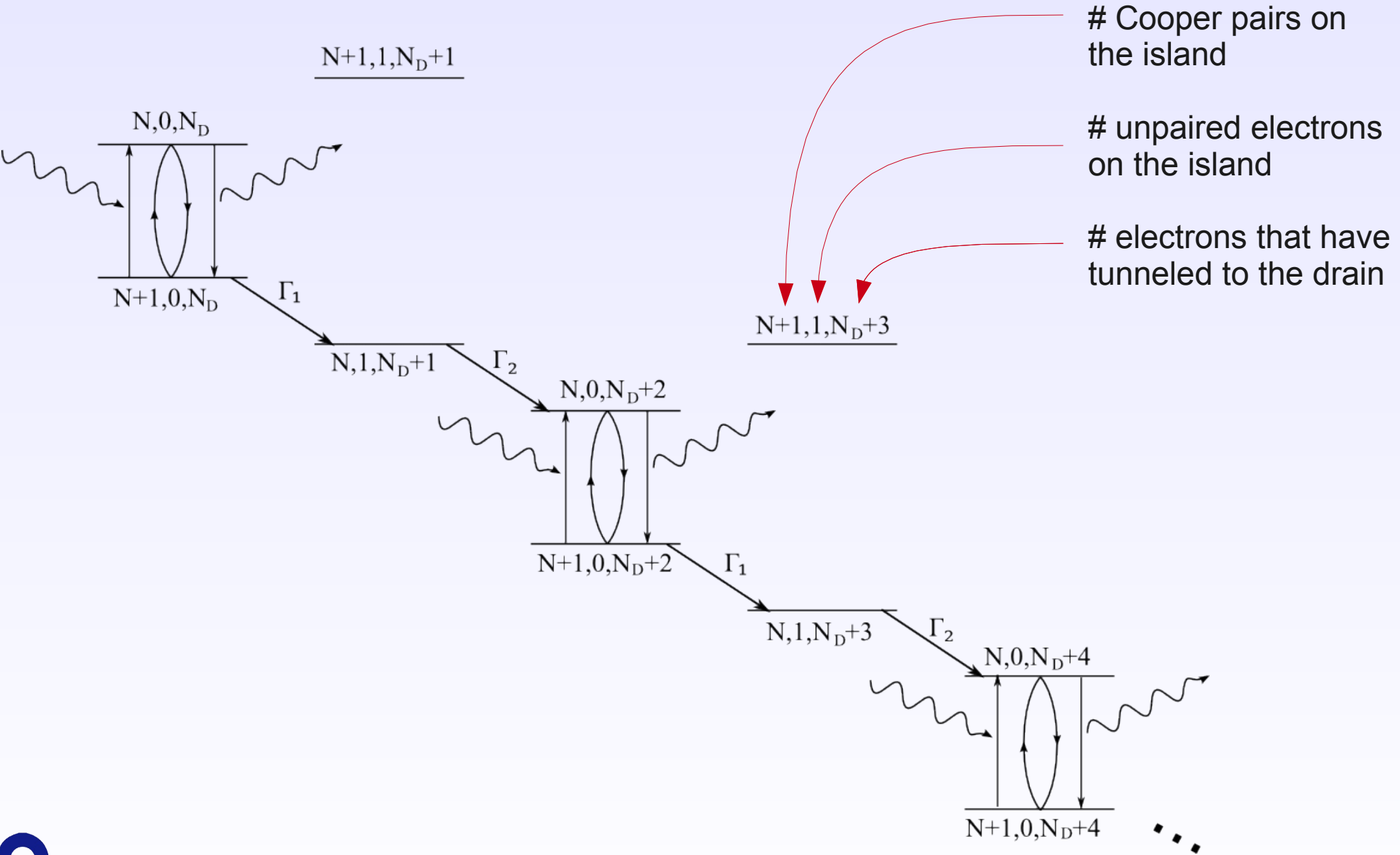


# Single-artificial-atom laser: pumping cycle

- *Population inversion of the artificial atom by a Josephson quasi-particle cycle:*



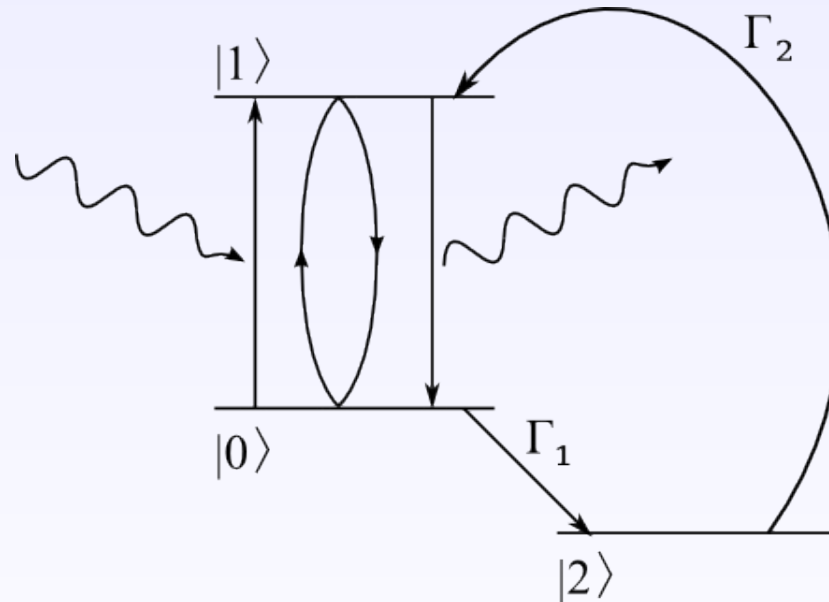
# Single-artificial-atom laser: lasing cascade





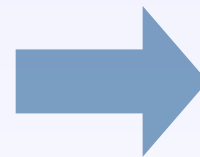
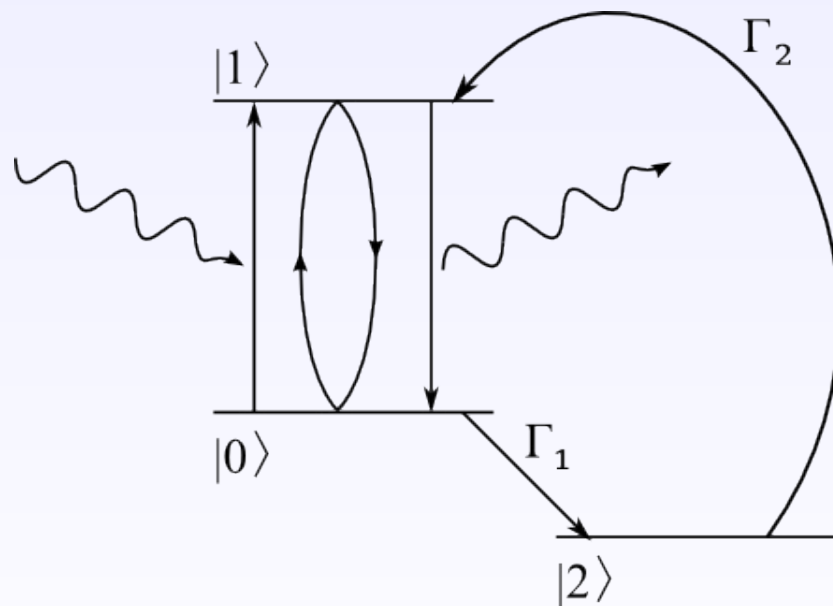
# Single-artificial-atom laser: lasing cycle

- *The artificial atom can be described as a three-level system*
- *Only two of the levels are coupled to the cavity mode*
- *Population inversion via a pumping cycle that involves the third state*

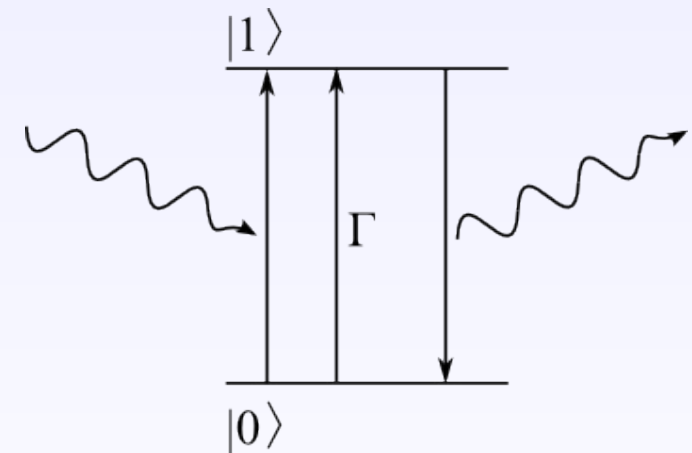


# Single-artificial-atom laser: simplified model

- *To begin with:*
  - *Only keep the two levels that couple to the cavity*
  - *Replace the pumping cycle with a reversed relaxation process*

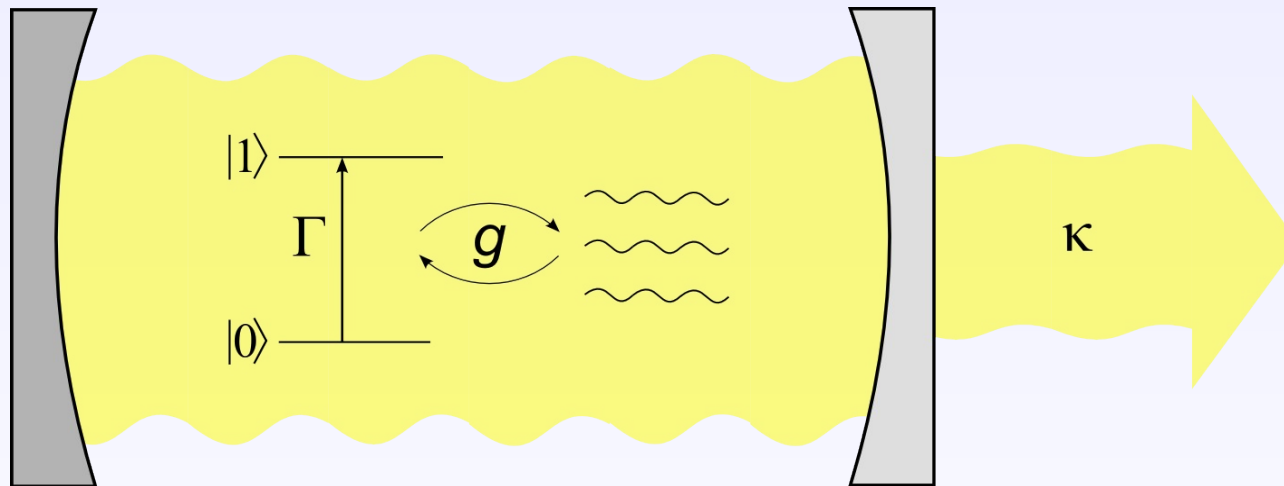


## Simplified atom



# A two-level atom in a cavity

- Atom pumping rate  $\Gamma$
- Atom/cavity interaction strength  $g$
- Cavity relaxation rate  $\kappa$



For lasing to occur:  $\Gamma, g \gg \kappa$

# A two-level atom in a cavity: Hamiltonian

- *Jaynes-Cummings model*

$$\hat{H} = \frac{\hbar \omega_a}{2} \hat{\sigma}_z + \hbar \omega_0 \hat{a}^\dagger \hat{a} + g \sigma_x (\hat{a} + \hat{a}^\dagger)$$

$\omega_a$  = natural frequency of atom

$\omega_0$  = natural frequency of cavity/resonator

$g$  = bare atom/cavity interaction strength

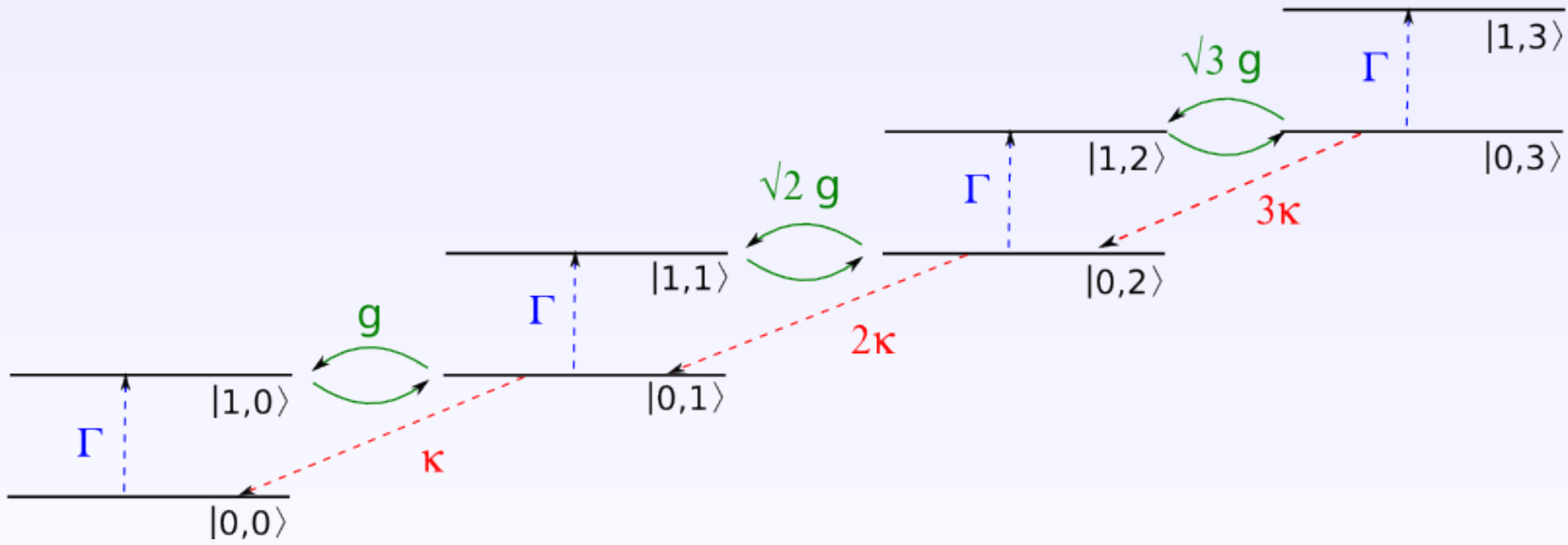
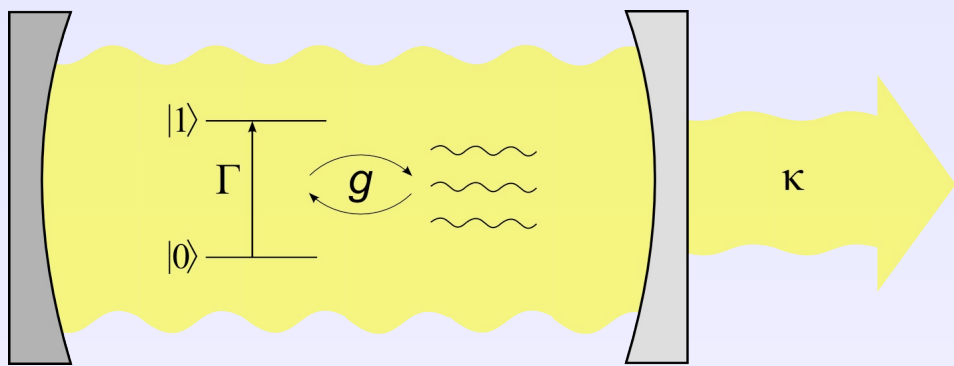
- *Atom/cavity in resonance*

$$\omega_a = \omega,$$

- *The state of the atom/cavity system is denoted:*

$$\left| n_a, n_c \right\rangle \Leftrightarrow \begin{array}{l} n_a = \# \text{ excitations in the atom (0 or 1)} \\ n_c = \# \text{ excitations in the cavity (0, 1, 2, \dots)} \end{array}$$

# Energy levels and relaxation/pumping rates



*For lasing:*  $\Gamma, g \gg \kappa$

# Rate equations: Cavity loss and atom emission

- *Cavity loss (~ laser output)*

$$\Gamma_{loss} = n \kappa$$

- *Atom emission (to cavity)*

- *Small n:  $\Gamma \gg \sqrt{n} g$*

$$\Gamma_{emission} = \frac{4 n g^2}{\Gamma}$$

- *Large n:  $\Gamma \ll \sqrt{n} g$*

$$\Gamma_{emission} \rightarrow \frac{\Gamma}{2}$$

*(maximal emission rate)*

# Rate equations: Cavity loss and atom emission

- Cavity loss ( $\sim$  laser output)

$$\Gamma_{loss} = n \kappa$$

- Atom emission (to cavity)

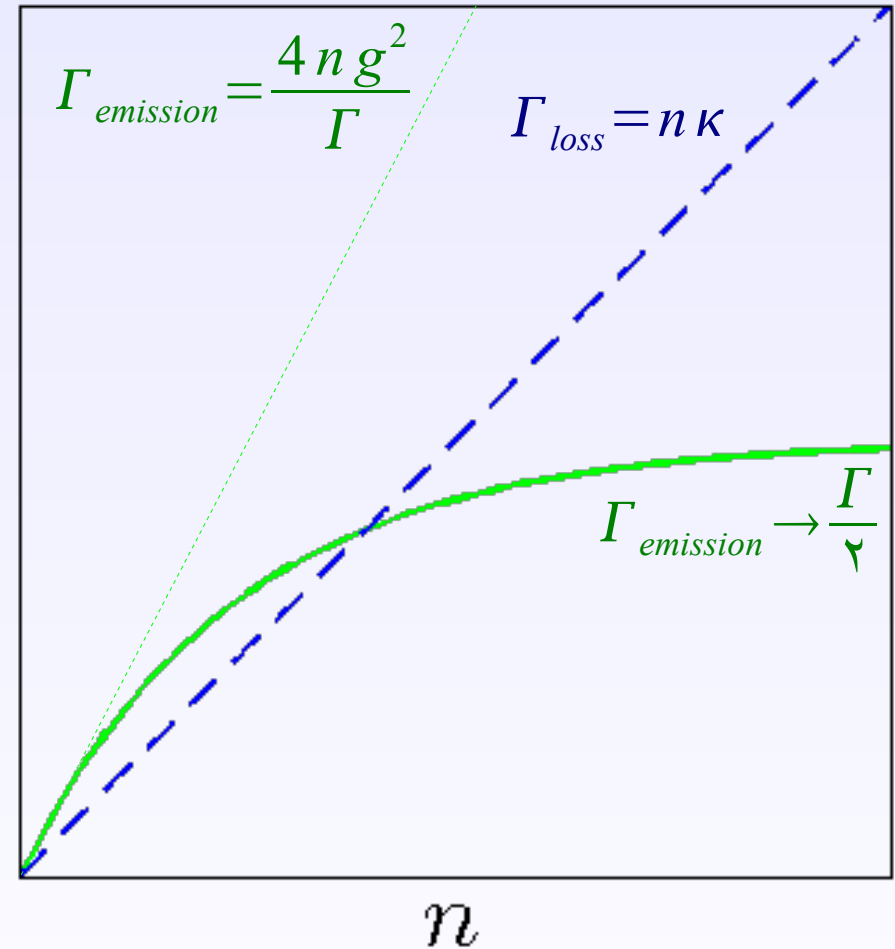
– Small  $n$ :  $\Gamma \gg \sqrt{n} g$

$$\Gamma_{emission} = \frac{4 n g^2}{\Gamma}$$

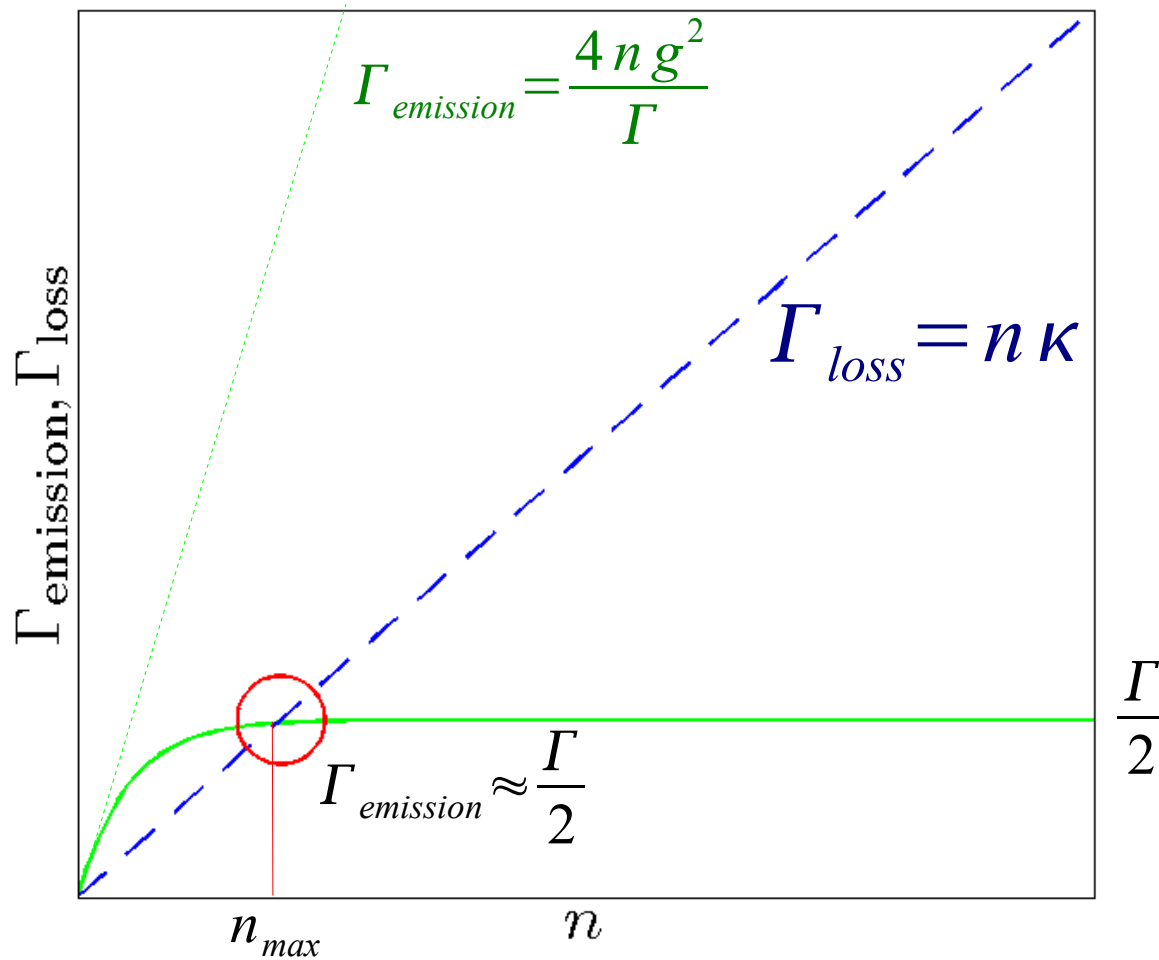
– Large  $n$ :  $\Gamma \ll \sqrt{n} g$

$$\Gamma_{emission} \rightarrow \frac{\Gamma}{2}$$

(maximal emission rate)



# Rate equations: Cavity loss and atom emission



State with maximal occupation probability is given by

Detailed balance:

$$\Gamma_{loss}(n) = \Gamma_{emission}(n)$$

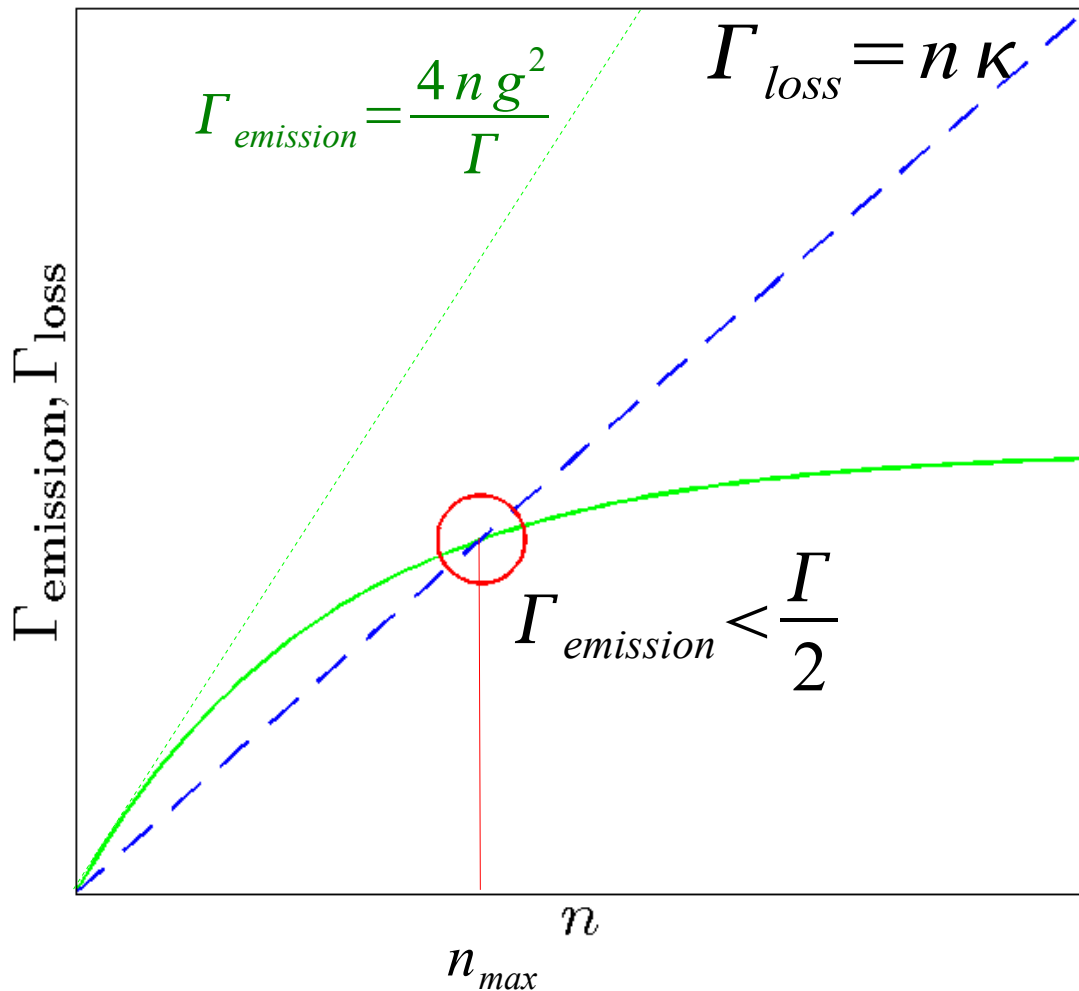


$$n_{max} \approx \frac{\Gamma}{\kappa}$$



# Rate equations: Cavity loss and atom emission

Increasing  $\Gamma$ :



State with maximal occupation probability is given by

Detailed balance:

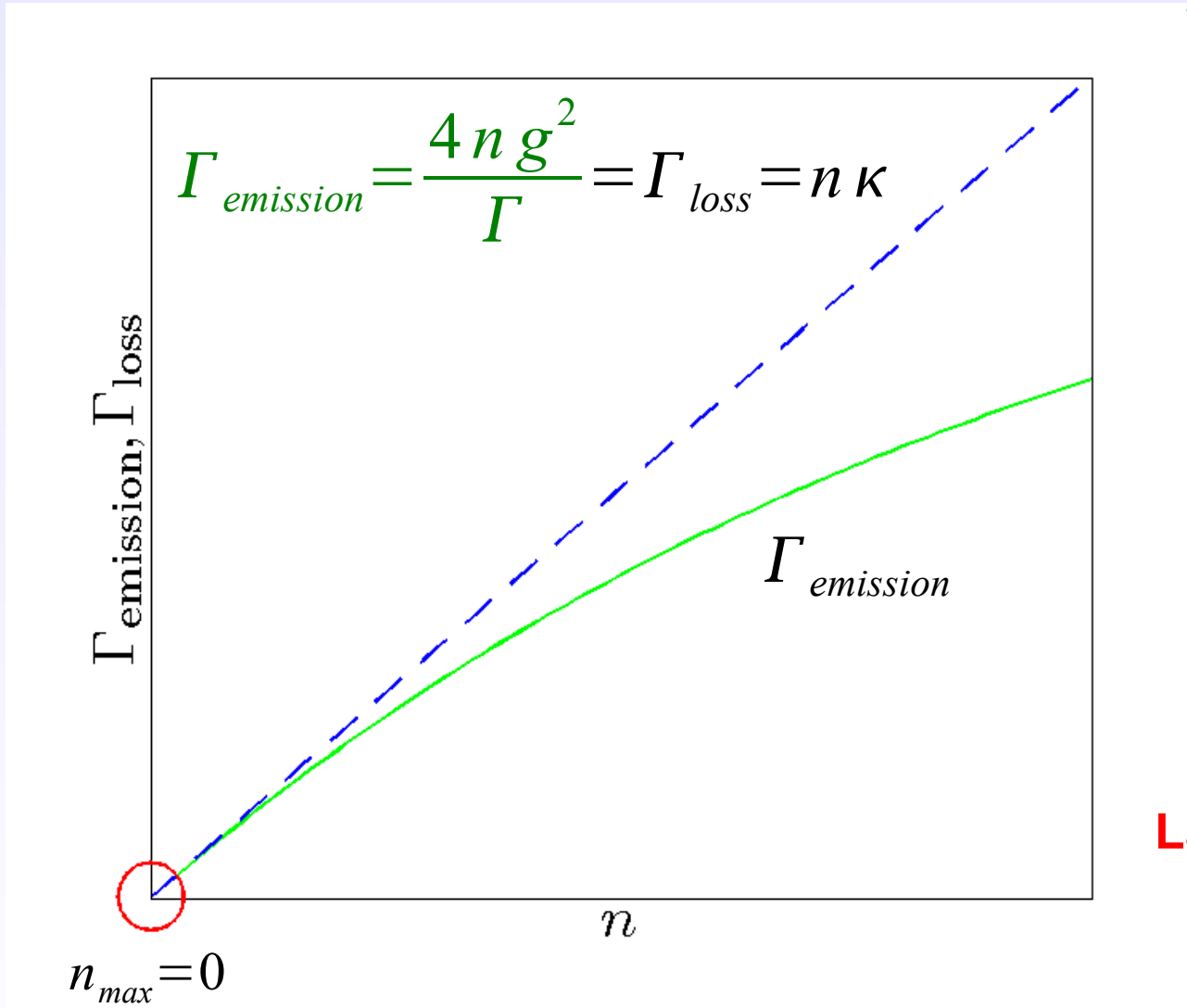
$$\Gamma_{\text{loss}}(n) = \Gamma_{\text{emission}}(n)$$



$$n_{\text{max}} < \frac{\Gamma}{2 \kappa}$$

# Rate equations: Lasing suppression threshold

Further increasing  $\Gamma$ :



When the cavity loss rate exceeds the atom emission rate:

no build-up in the cavity

$$\Gamma_{loss}(n) > \Gamma_{emission}(n)$$



**Lasing suppression threshold**

$$\frac{\xi g^2}{\Gamma} < \kappa$$

# Semi-classical and numerical calculations

## Lindblad master equation

$$\frac{d\rho}{dt} = -\frac{i}{\hbar} [\hat{H}, \rho] + \Gamma \left( \hat{\sigma}_+ \rho \hat{\sigma}_- - \frac{1}{2} \hat{\sigma}_- \hat{\sigma}_+ \rho - \frac{1}{2} \rho \hat{\sigma}_- \hat{\sigma}_+ \right) + \kappa \left( a \rho a^\dagger - \frac{1}{2} a^\dagger a \rho - \frac{1}{2} \rho a^\dagger a \right)$$

$\kappa$  is the *relaxation* rate of the oscillator

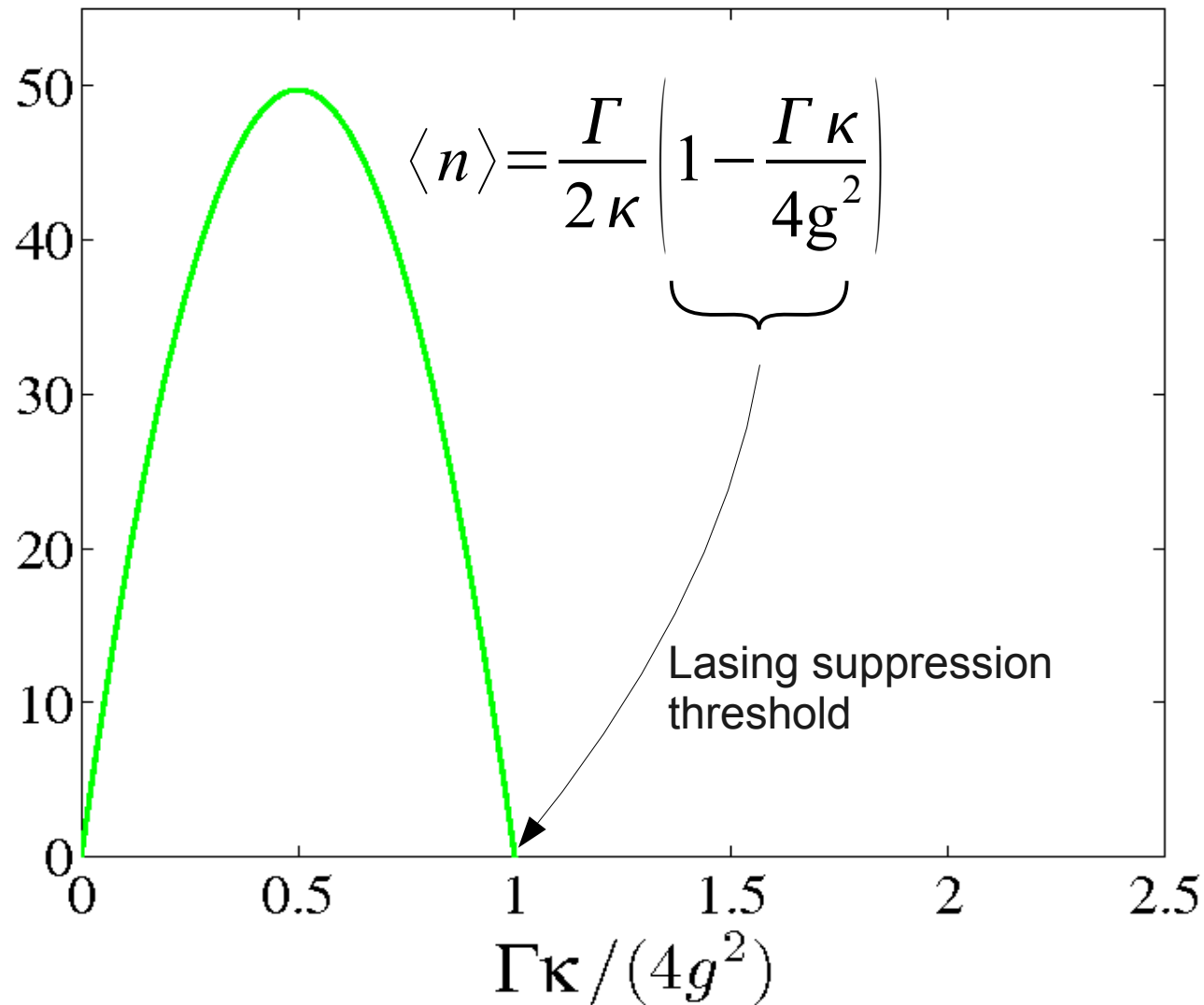
$\Gamma$  is the *excitation* rate of the two-level system

*Using mean-field approximation we can derive the average photon number in the cavity (for the steady state)*

$$\langle n \rangle = \frac{\Gamma}{2\kappa} \left( 1 - \frac{\Gamma\kappa}{4g^2} \right)$$

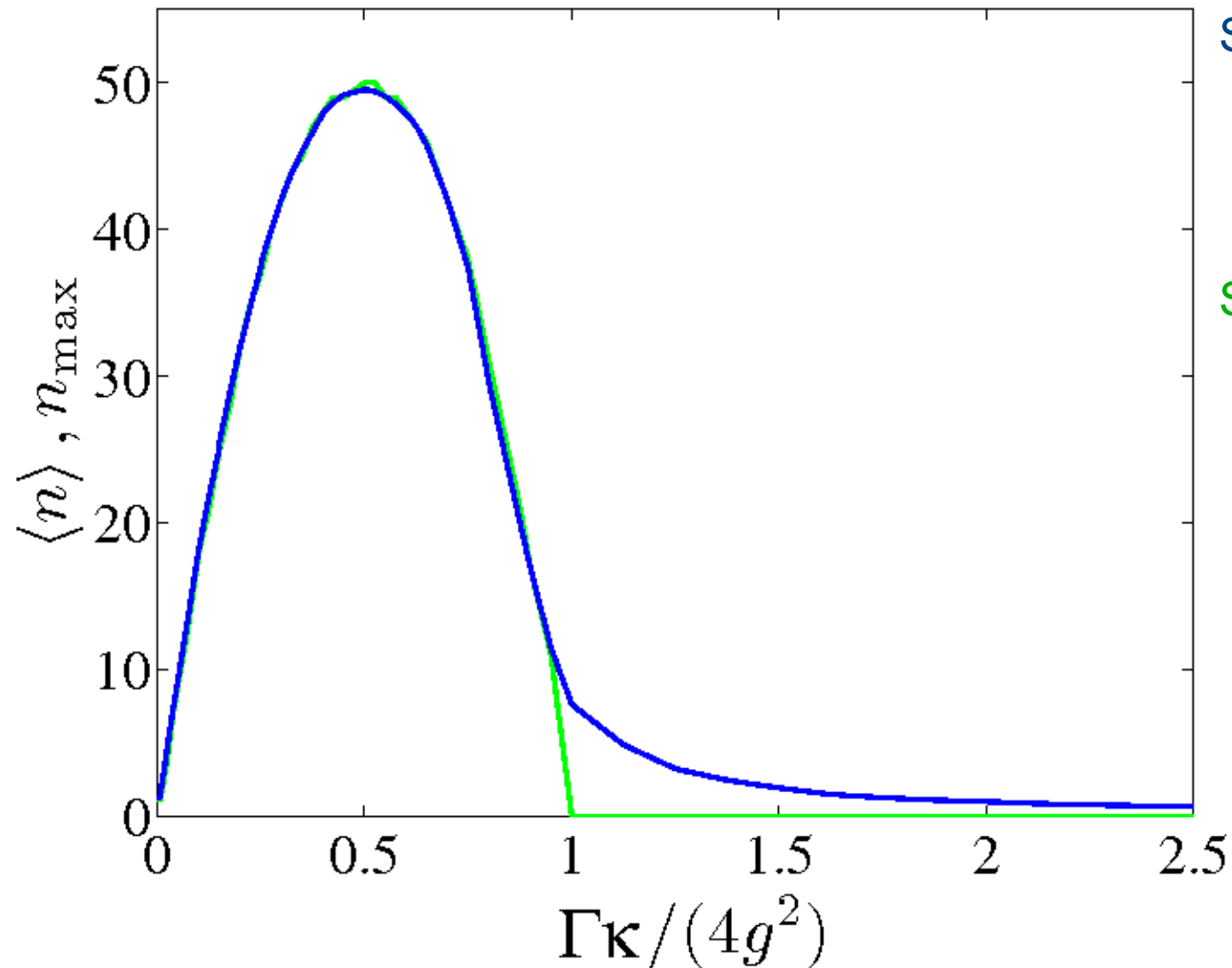
# Semi-classical results

- Average photon occupation number calculated with the *semi-classical equations of motion*



# Numerical results

- Average photon occupation number calculated by *numerical integration of the master equation*

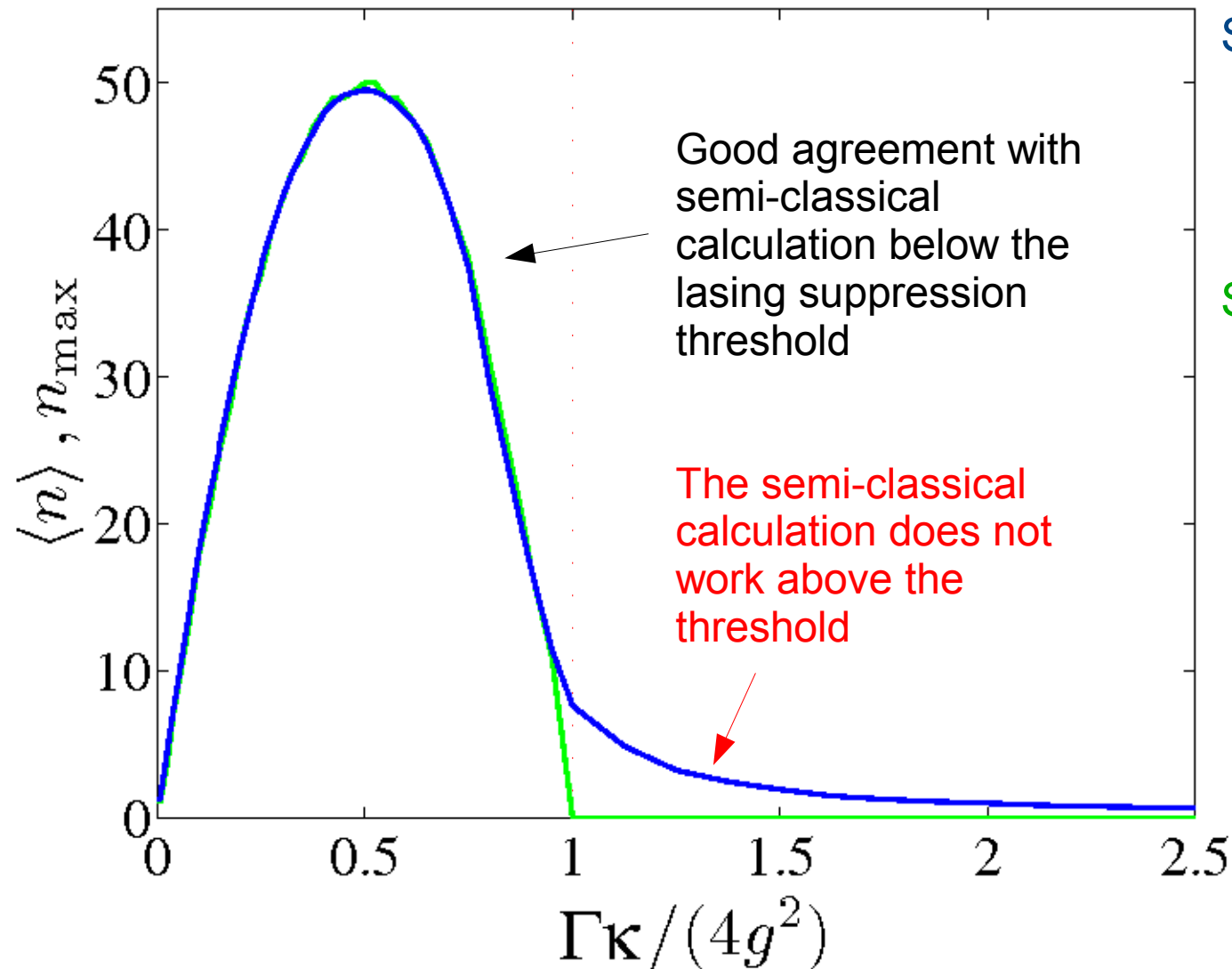


Solid blue curve:  
Average photon  
occupation  
number

Solid green curve:  
State with  
maximum  
occupation

# Comparison: semi-classical and numerical results

- Good agreement between semi-classical and numerical results below the lasing suppression threshold

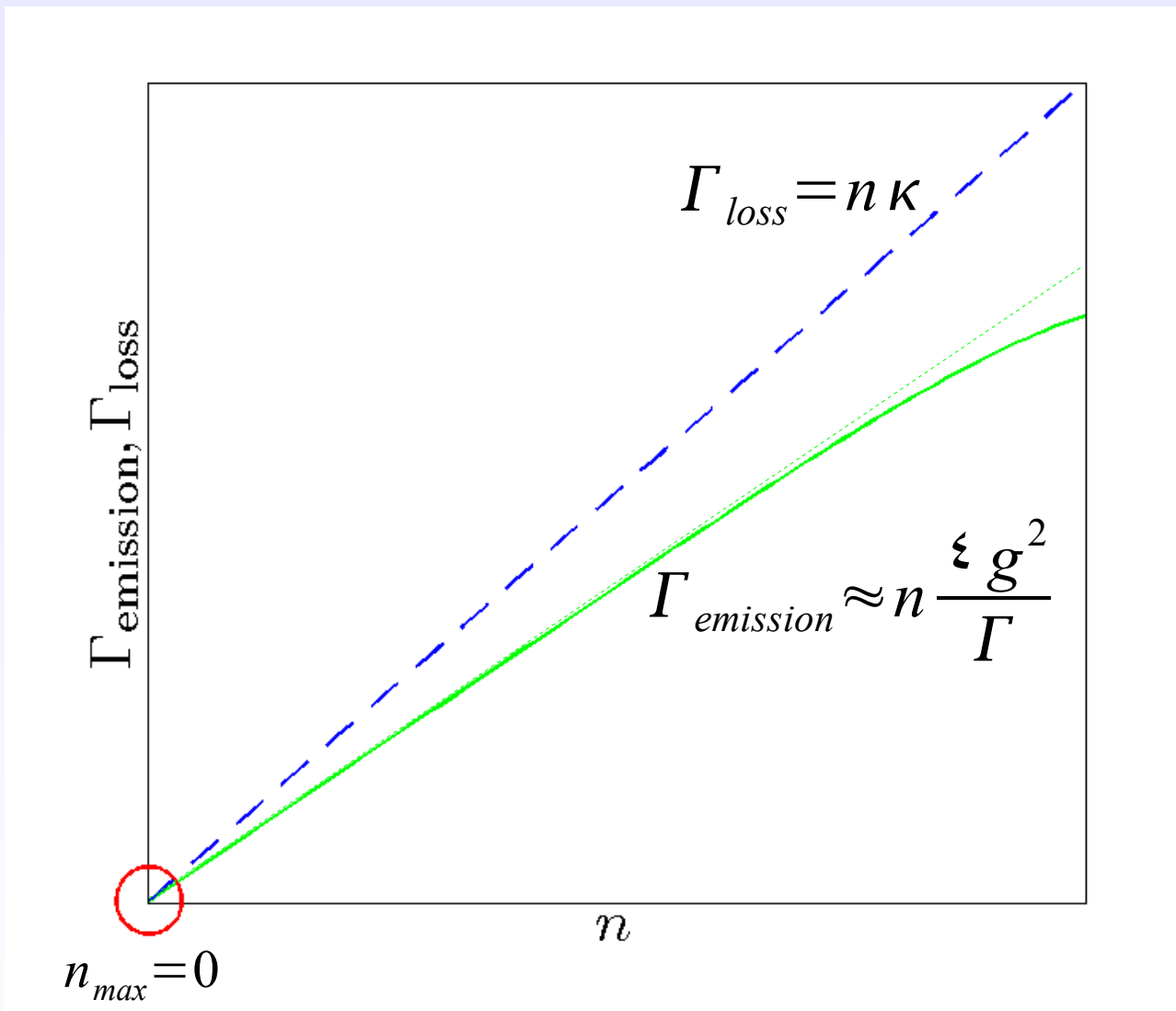


Solid blue curve:  
Average photon  
occupation  
number

Solid green curve:  
State with  
maximum  
occupation

# State of the cavity above the threshold

Going back to rate equations:



Detailed balance:

$$\frac{P_{n+1}}{P_n} \approx \frac{\Gamma_{\text{emission}}}{\Gamma_{\text{loss}}}$$



$$\frac{P_{n+1}}{P_n} \approx \frac{4g^2}{\Gamma\kappa} = e^{\frac{\hbar\omega}{k_B T_{\text{eff}}}}$$

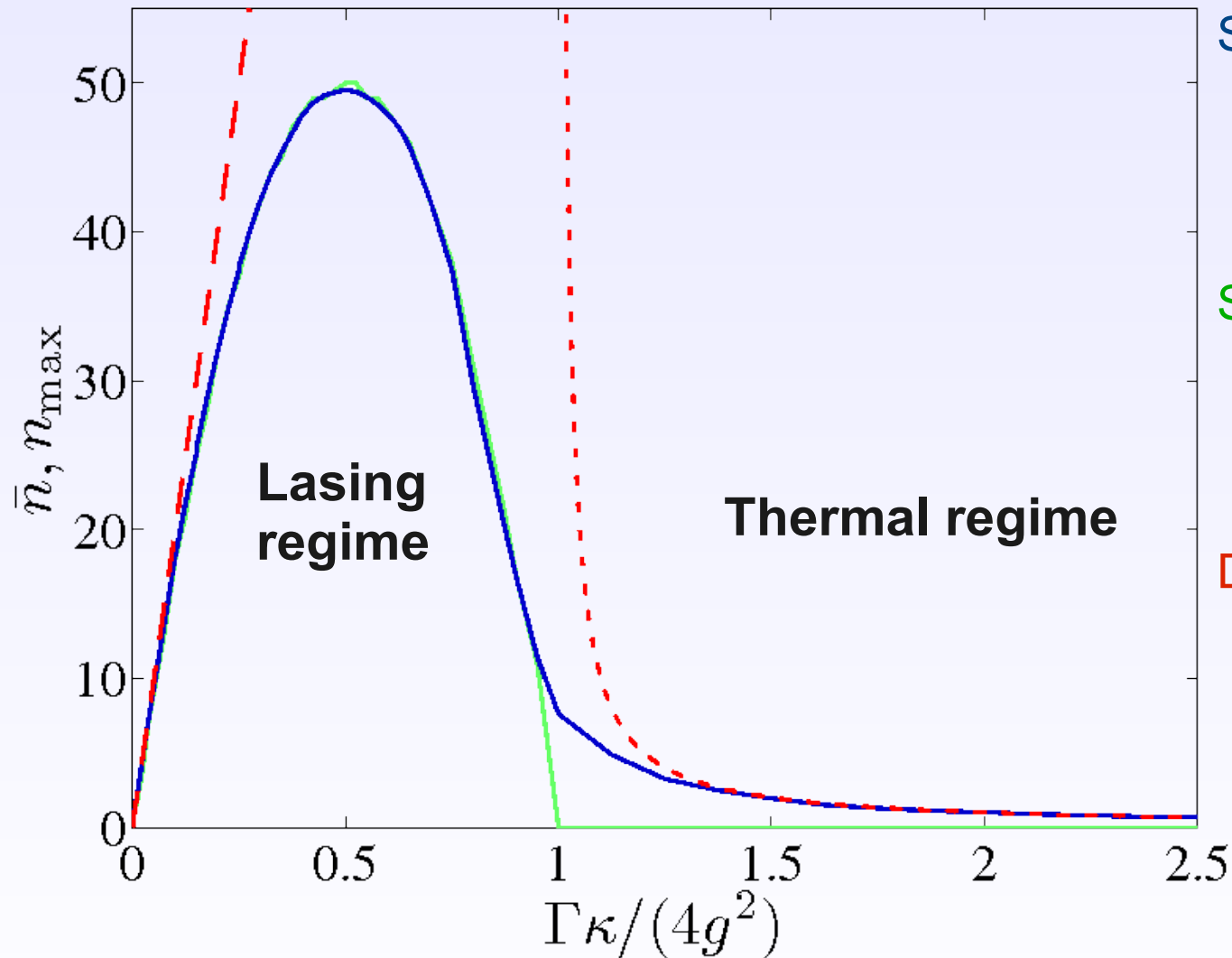


**Effective thermal-equilibrium state**

$$\langle n \rangle = \left( \frac{\Gamma\kappa}{4g^2} - 1 \right)^{-1}$$

# Analytical result for the thermal regime

- The analytical results for the regime above the lasing suppression threshold agree well the numerical calculations



Solid blue curve:  
Average photon  
occupation  
number

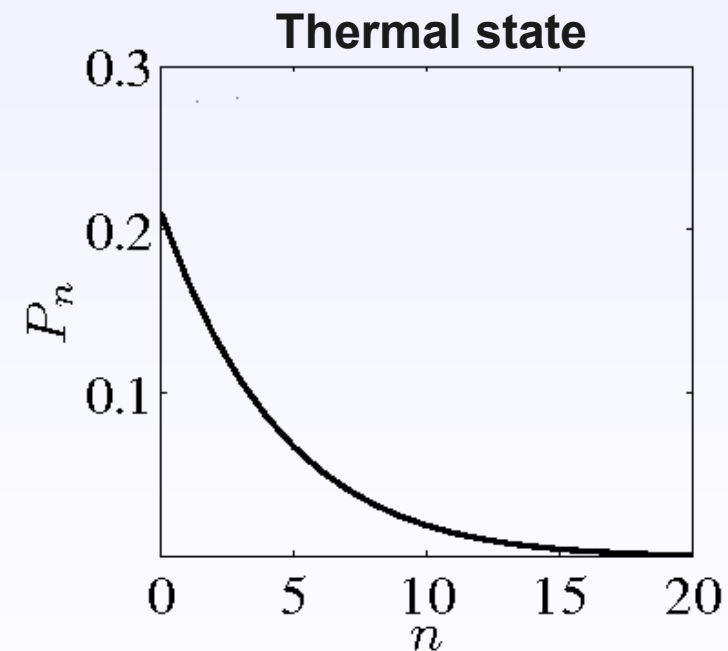
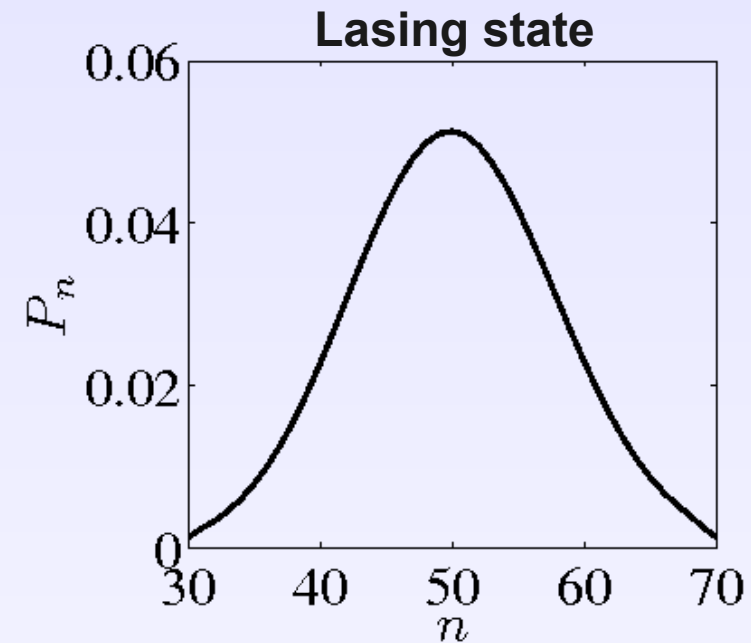
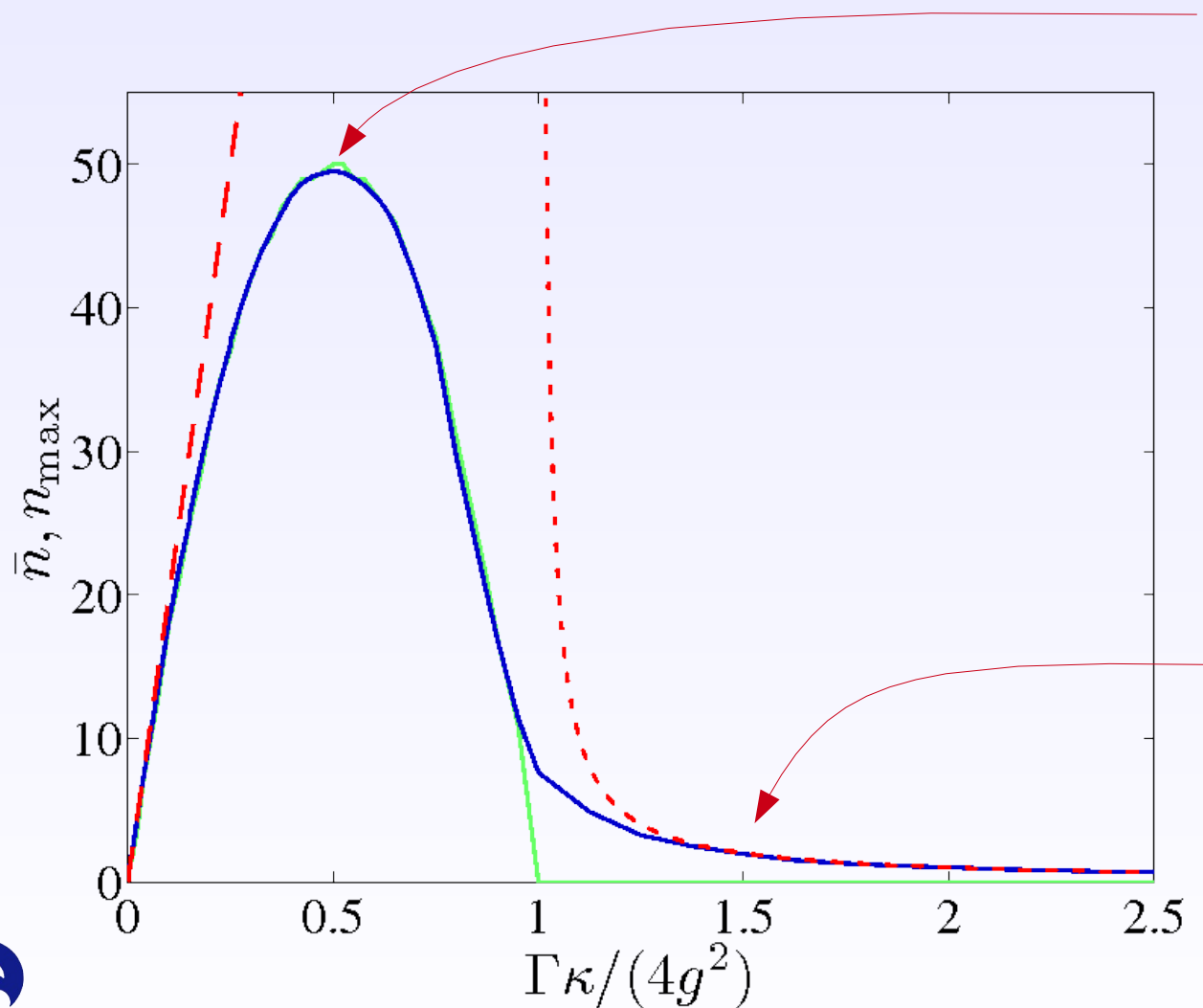
Solid green curve:  
State with  
maximum  
occupation

Dotted red curve  
Analytical result,  
thermal regime



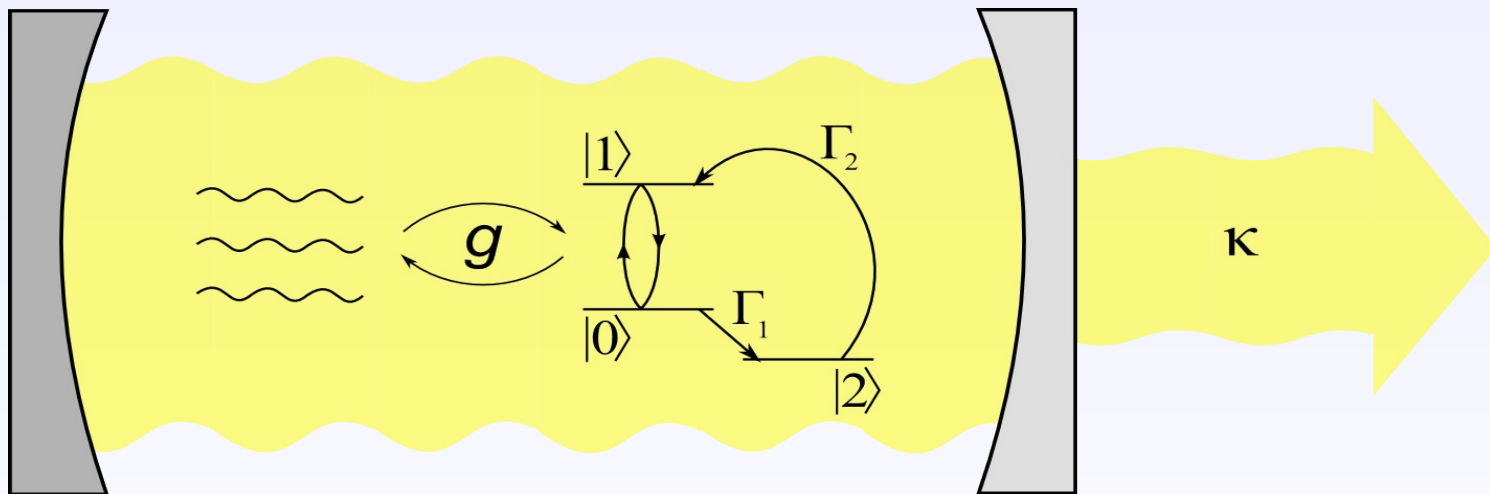
# Cavity photon-number distribution

- Representative cavity photon-number distributions for the lasing and the thermal regime.



# Three-level atom in a cavity

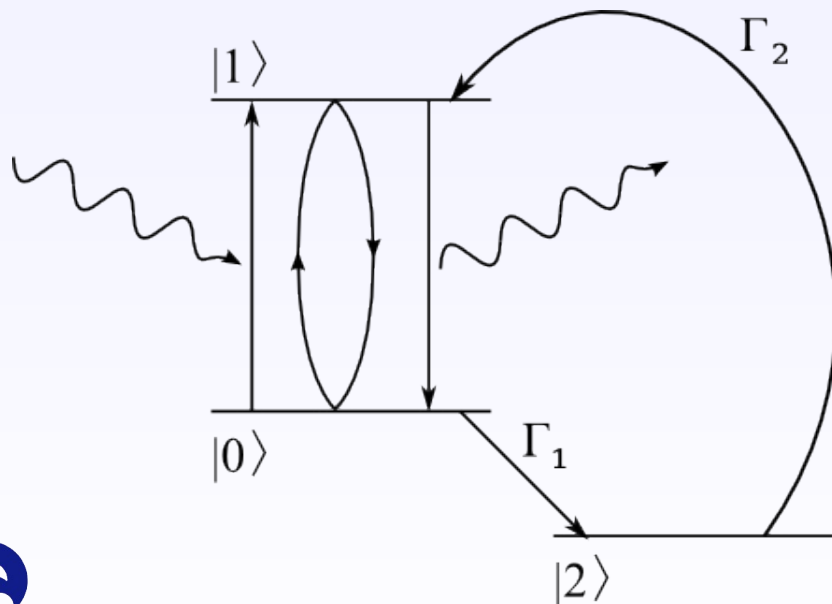
- Atom pumping rates  $\Gamma_1, \Gamma_2$
- Atom/cavity interaction strength  $g$
- Cavity relaxation rate  $\kappa$



# Three-level atom in a cavity: Hamiltonian

- Atom pumping rates  $\Gamma_1, \Gamma_2$
- Atom/cavity interaction strength  $g$
- Cavity relaxation rate  $\kappa$
- Jaynes-Cummings Hamiltonian

$$\hat{H} = \frac{\hbar \omega_a}{2} (\sin \theta \hat{\sigma}_z + \cos \theta \hat{\sigma}_x) + \hbar \omega_0 \hat{a}^\dagger \hat{a} + g_0 \sigma_z (\hat{a} + \hat{a}^\dagger)$$



The Pauli matrices acts on the atomic states:

$$|0\rangle, |1\rangle$$

# Three-level atom in a cavity: Analytical results

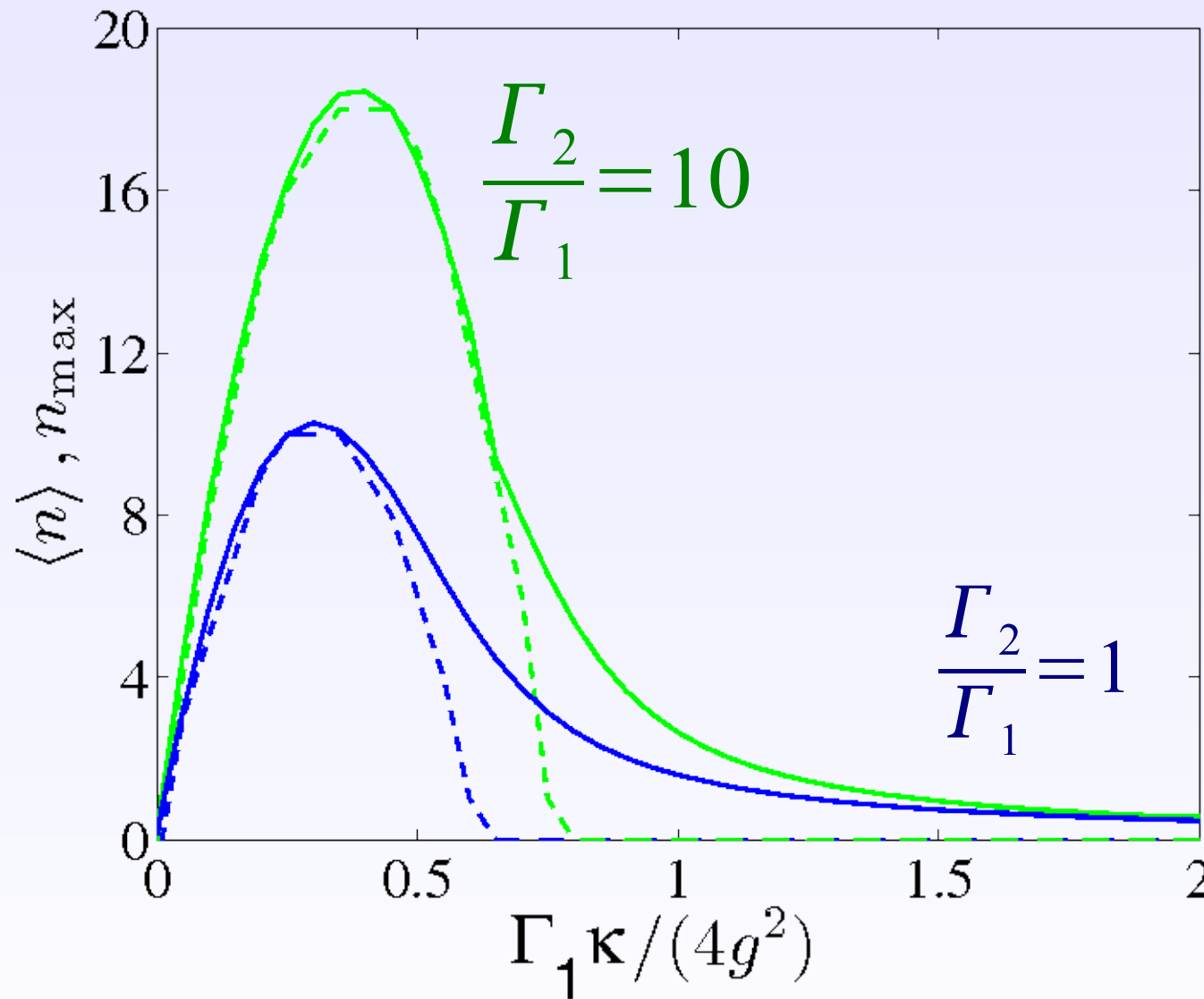
- *Lasing threshold:*

$$\frac{\Gamma_1 \kappa}{4g^2} = \frac{\cos \theta}{\cos^2 \theta + \left( \frac{1}{2} + \frac{\Gamma_1}{4\Gamma_2} \right) \sin^2 \theta}$$

- *Average photon occupation number in the cavity from semi-classical calculation*

$$\langle n \rangle = \frac{\Gamma_1}{2\kappa} \left( \frac{1}{1 + \frac{\Gamma_1}{2\Gamma_2}} \cos \theta - \left( 1 + \frac{1 - \frac{\Gamma_1}{2\Gamma_2}}{1 + \frac{\Gamma_1}{2\Gamma_2}} \cos^2 \theta \right) \frac{\Gamma_1^2 \kappa}{8g^2} \right)$$

# Numerical results for three-level-atom model



Solid curves:  
Average photon  
occupation  
number

Dashed curves:  
State with  
maximum  
occupation

# Conclusions

- *We study two models for single-atom lasing*
  - *Two-level atom*
  - *Three-level atom*
- *We analyzed these models using*
  - *Transition rate equations*
  - *Semi-classical equation of motion*
  - *Numerical simulations*
- *We found conditions for lasing*
- *We calculate the cavity photon-distribution*
  - *in lasing regime*
  - *in the suppressed-lasing regime*